Nuclear Dynamics at Extreme Conditions

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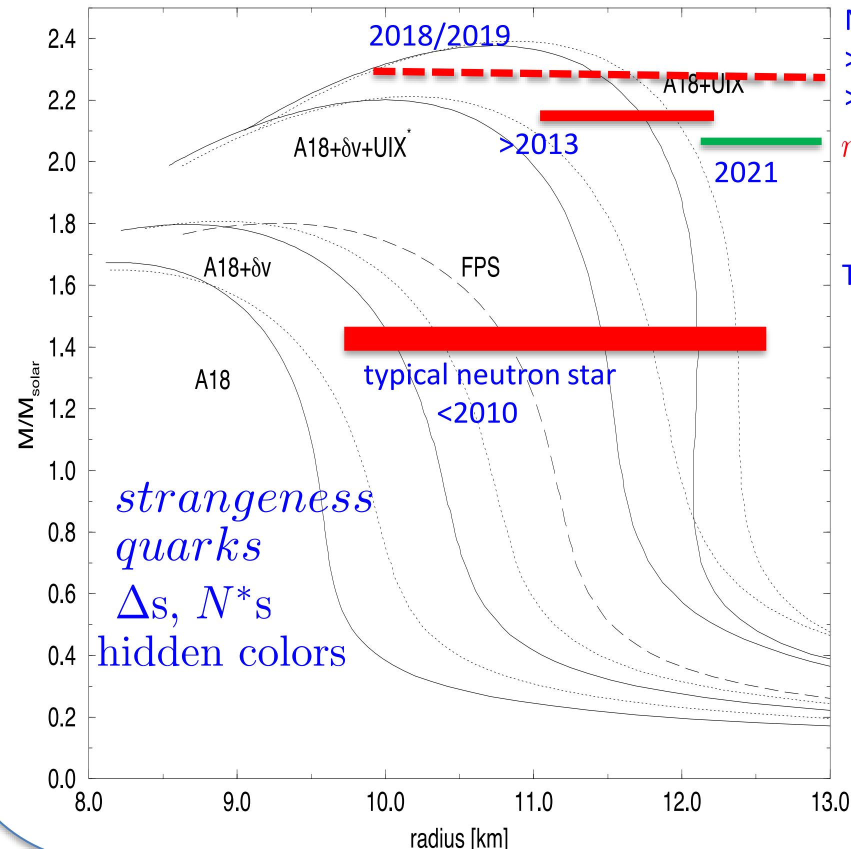
“Unreasonable” Persistence of Nucleons



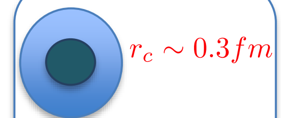
Neutron star masses
>2 solar masses
>2010

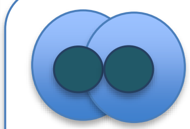
$r_{NN} \sim 0.6 - 0.8 fm$

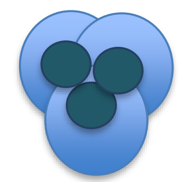
Typical neutron star
1.4 Solar mass

$r_{NN} \sim 0.7 - 1.0 fm$

H. Heiselberg,
V.Pandharipande
ARNPS 2000







three color singlet

“baryons”

Nuclear Dynamics at Short Distances

Probing NN and NNN Interactions at < 1fm

their role in the dynamics and structure disappearance in High Density Nuclear Matter

For NN interactions

- *Identification of NN interaction in the nuclear dynamics*–*PRC1993,2003, PRL2004- Nuclear Scaling*
- *Intermediate – short distance tensor forces* –*PRL 2006, Science2006 pn SRC dominance*
- *Isospin dependence of the tensor forces, momentum sharing*–*PRC14, Science2014, Nature 2018*
  - *NN repulsive core*
  - *Hadron-quark transition in the core* *PRL2023*
  - *non-nucleonic components, hidden color, gluons*

For NNN interactions

- *Identification of NNN interaction in the nuclear dynamics* *PRC2019, 2023*
- *Evaluation of irreducible 3N forces*
- *Evaluation of non-nucleonic component in NNN interactions*

Probing NN Repulsive Core

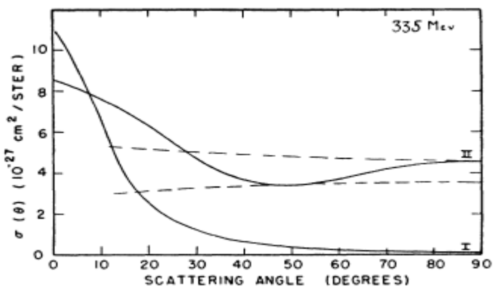
- NN force is attractive: But Nuclei are Stable

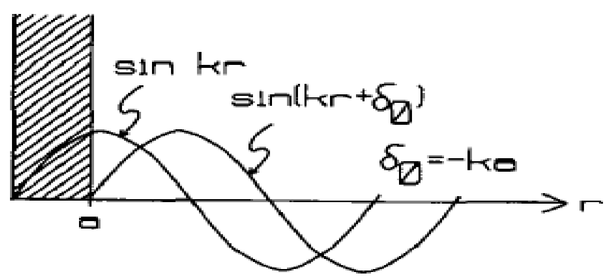
"If the two-body forces are everywhere attractive and if many-body forces are neglected then the nucleon pairs are sufficiently close to take advantage of attractive interactions and a collapsed state of nuclear matter results "

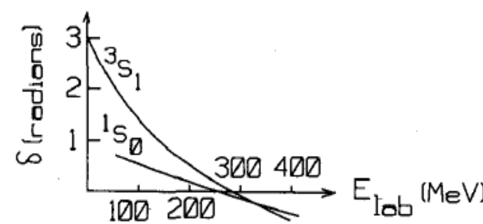
G. Breit and E.P. Wigner, Phys. Rev. 53, 998 (1938).

Many body forces keeping nucleus stable

Jastrow 1951 assumed the existence of the infinite hard core to explain the angular distribution of pp cross section at 340 MeV (r_0=0.6fm)







Non-monotonic NN central potential with the repulsive core was introduced: Brueckner & Watson 1953 to obtain nuclear density saturation.

Modern NN Potentials

$$V^{2N} = V_{EM}^{2N} + V_{\pi}^{2N} + V_{R}^{2N}$$

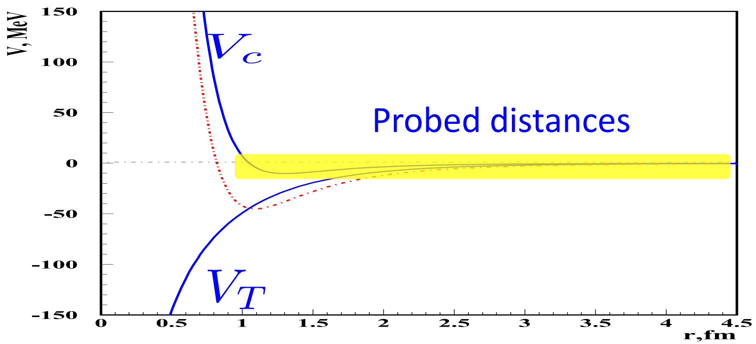
$$V_R^{2N} = V^c + V^{l2} L^2 + V^t S_{12} + V^{ls} L \cdot S + v^{ls2} (L \cdot S)^2$$

$$V^i = V_{int,R} + V_{core}$$

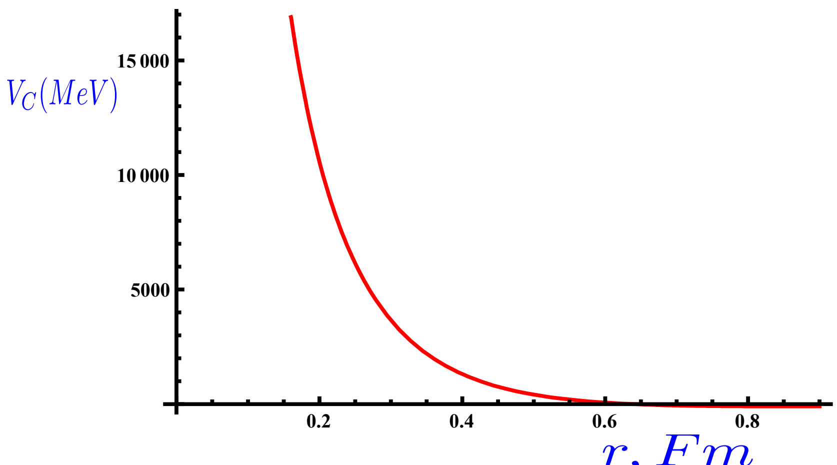
$$V_{core} = \left[1 + e^{\frac{r - r_0}{a}} \right]^{-1}$$

60’s

Currently: Probed NN structure up to > 0.8fm

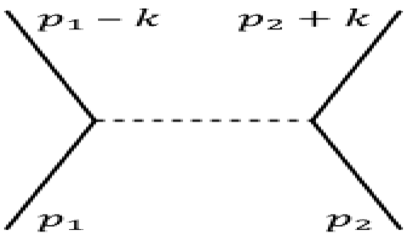


Next: NN – Repulsive Core



Nuclear Forces and Field Theory

- 1935 meson exchange model: H. Yukawa



$$V(r) = -g^2 \frac{e^{-mr}}{r}.$$

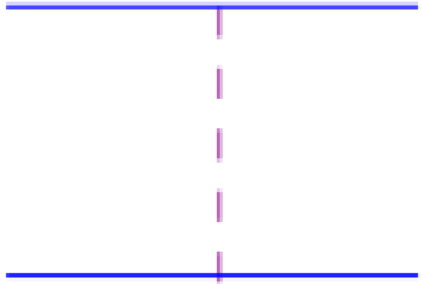
- Predicted meson with around 100 MeV :

- 1947 meson discovered: C.F. Powel

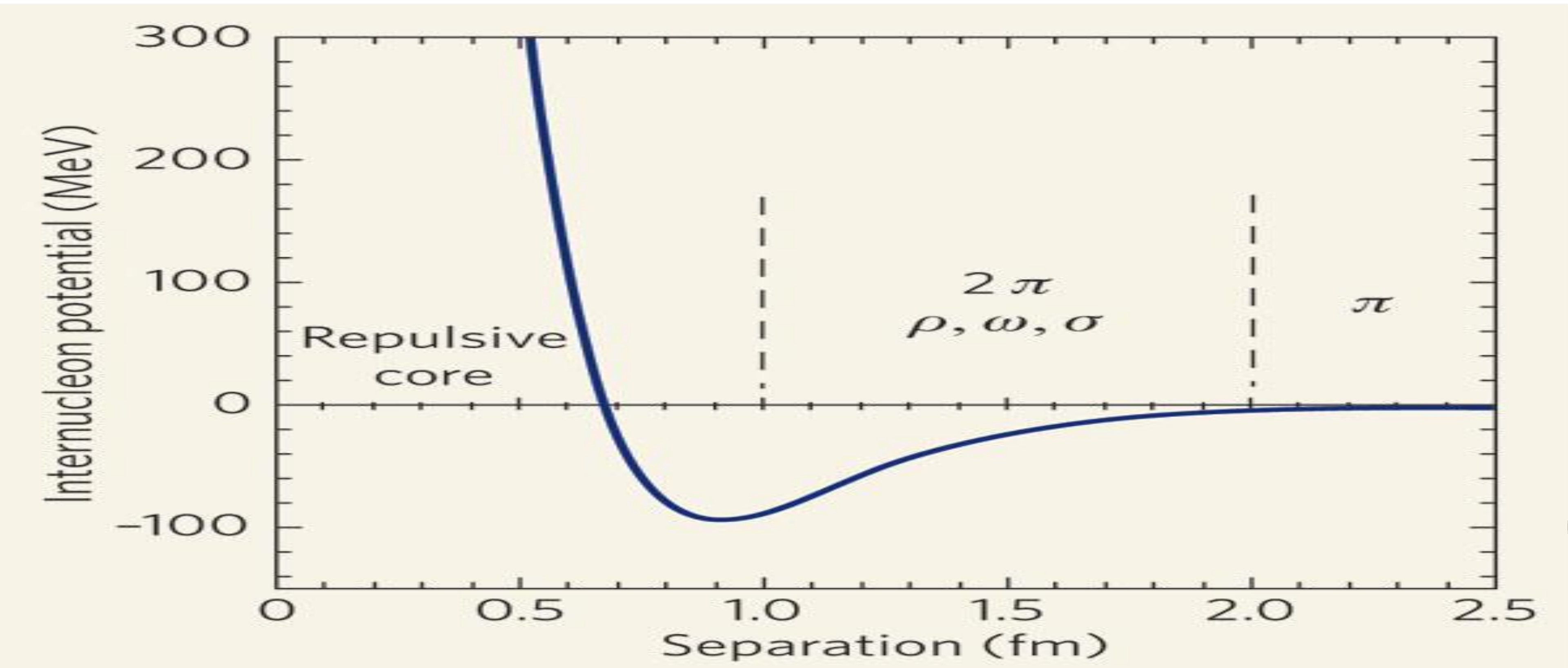
- 1943-1945 – seen by:

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

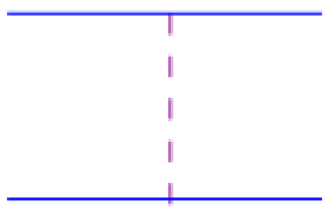
Artem Alikhanian
Aragats Cosmic Ray Station



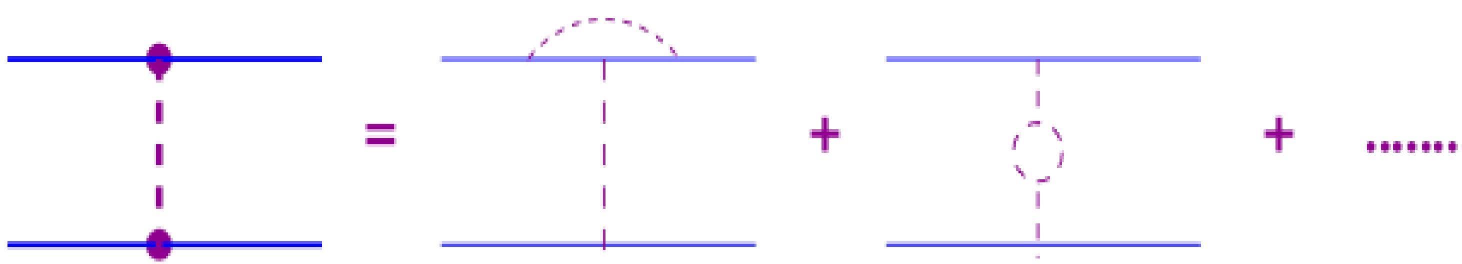
$\sigma, \quad \pi, \quad \rho, \quad \omega, \ldots$



Field Theory of Nucleons & Mesons 1947-



$\sigma, \quad \pi, \quad \rho, \quad \omega, \ldots$



Pomeranchuk, Landau - 1950's

$$g^2(\Lambda^2) = \frac{g^2}{1 - 5(\frac{g^2}{4\pi})ln(\frac{\Lambda^2}{m^2})}$$

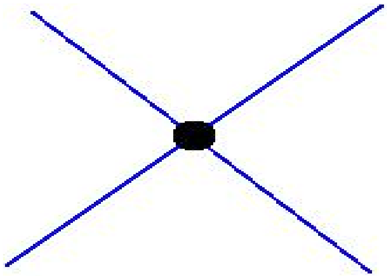
$$\alpha_{em}(Q^2) = \frac{\alpha_{em}(\mu^2)}{1 - \frac{\alpha_{em}(\mu^2)}{3\pi} ln\frac{Q^2}{\mu^2}}$$

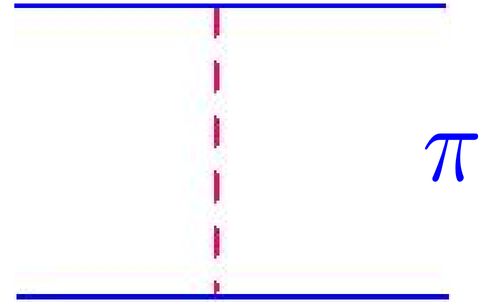
$$Q = m_e e^{\frac{3\pi}{2\alpha}} \sim 10^{277} GeV$$

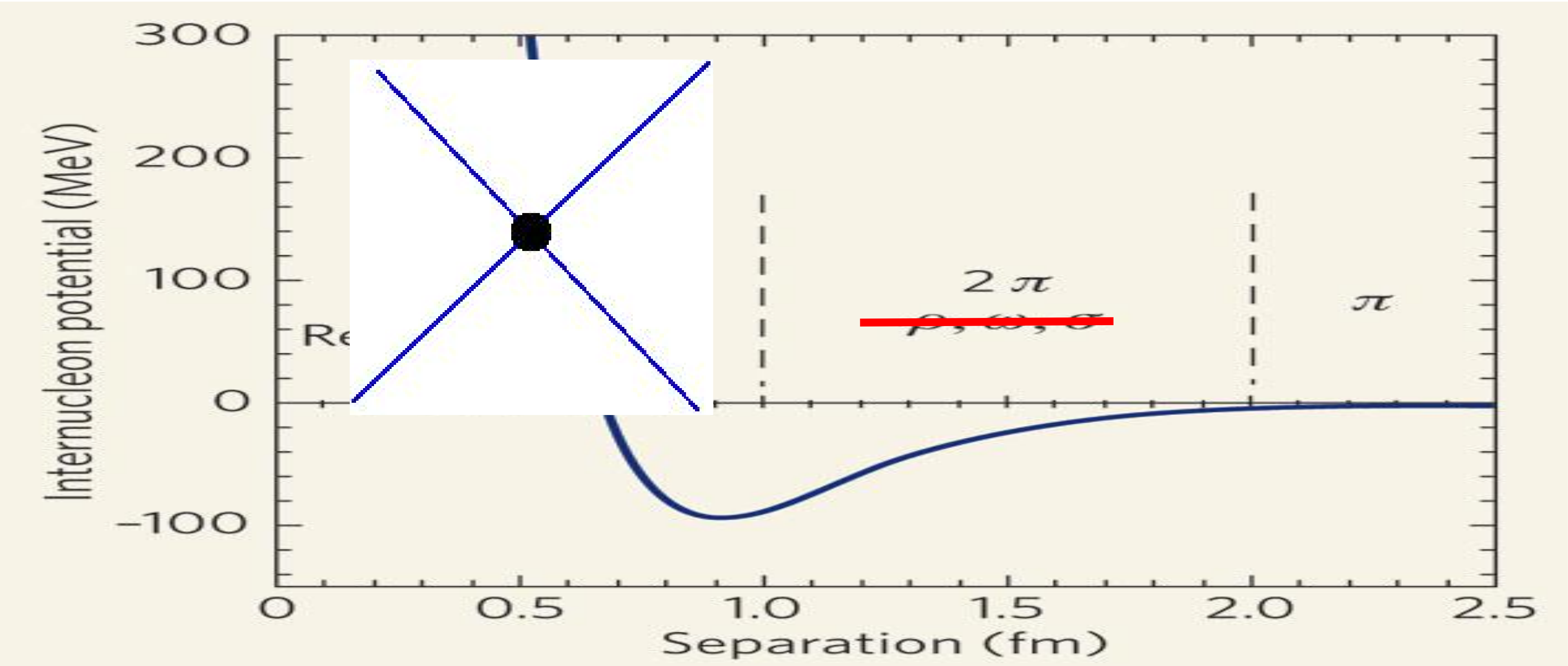
Infinite interaction occurs at transferred momenta approx 500 MeV/c or at internucleon distances 1~Fm.

It seems we have a problem about which the Nature is not aware of Y.Pomeranchuk

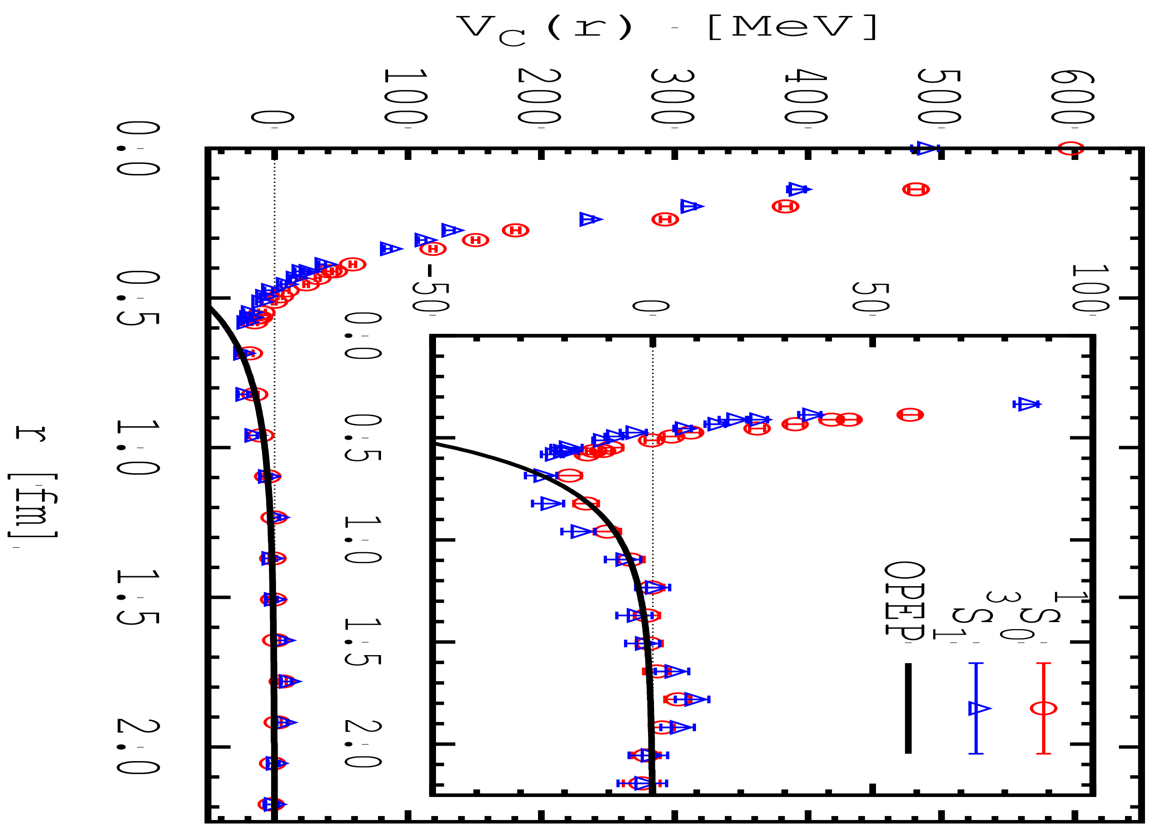
All formal quantum field theories with Yukawa type interactions contain the problem of the ``Zero Charge''
Pomeranchuk, Sudakov, Ter-Martirosyan, Phys. Rev. 1956



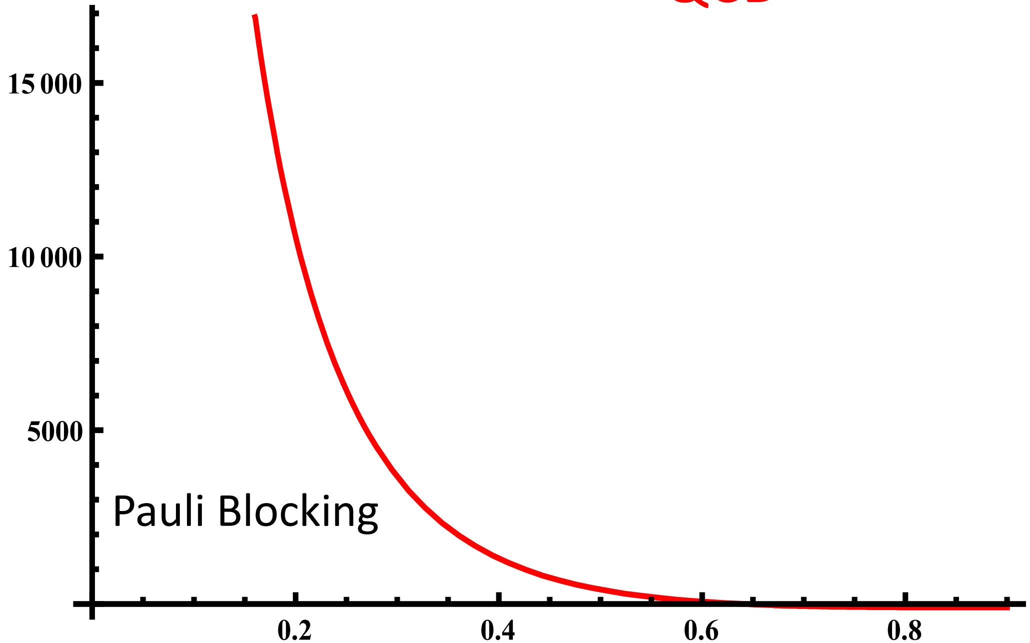

π


Internucleon potential (MeV)
300
200
100
0
−100
Re
2 π
ρ, ω, σ
π
0
0.5
1.0
1.5
2.0
2.5
Separation (fm)

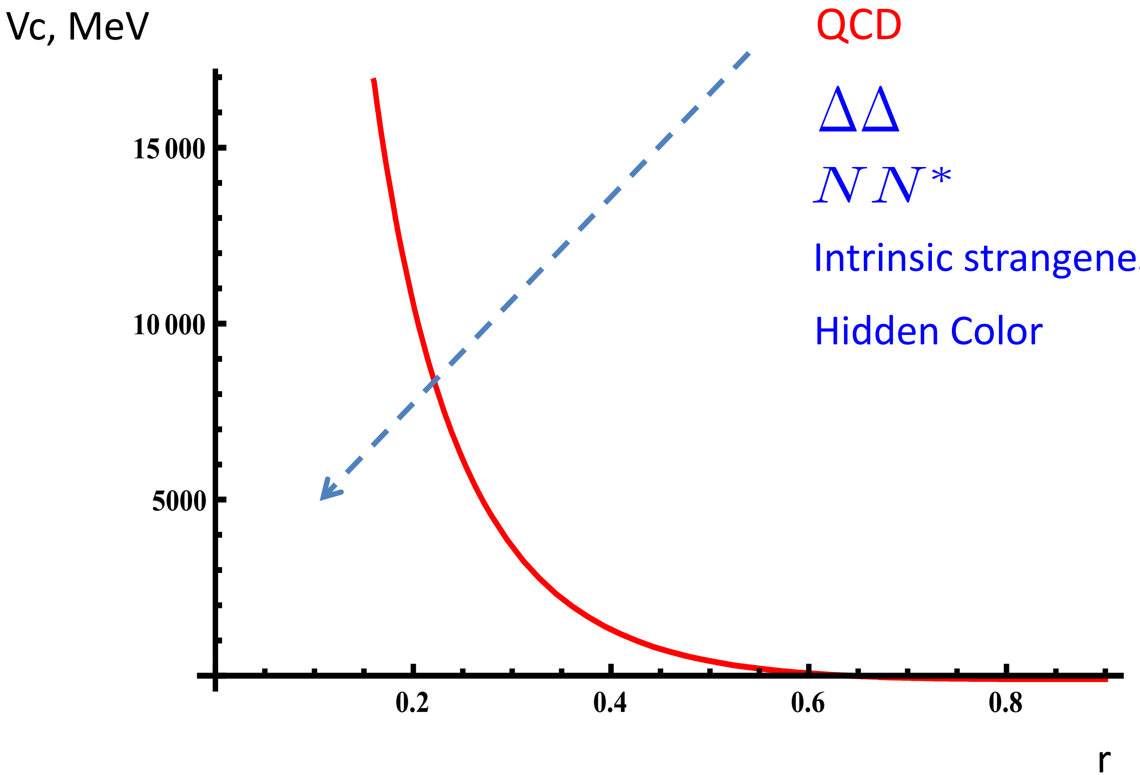
Lattice Calculations



QCD



Contradicts Neutron Star Observations:
will predict masses not more than 0.1 - 0.6 Solar mass



QCD

$\Delta\Delta$

$N N^*$

Intrinsic strangeness/charm

Hidden Color

~80% hidden color

Brodsky,Ji, Lepage, PRL 83

Probing the Deuteron at Short Distances

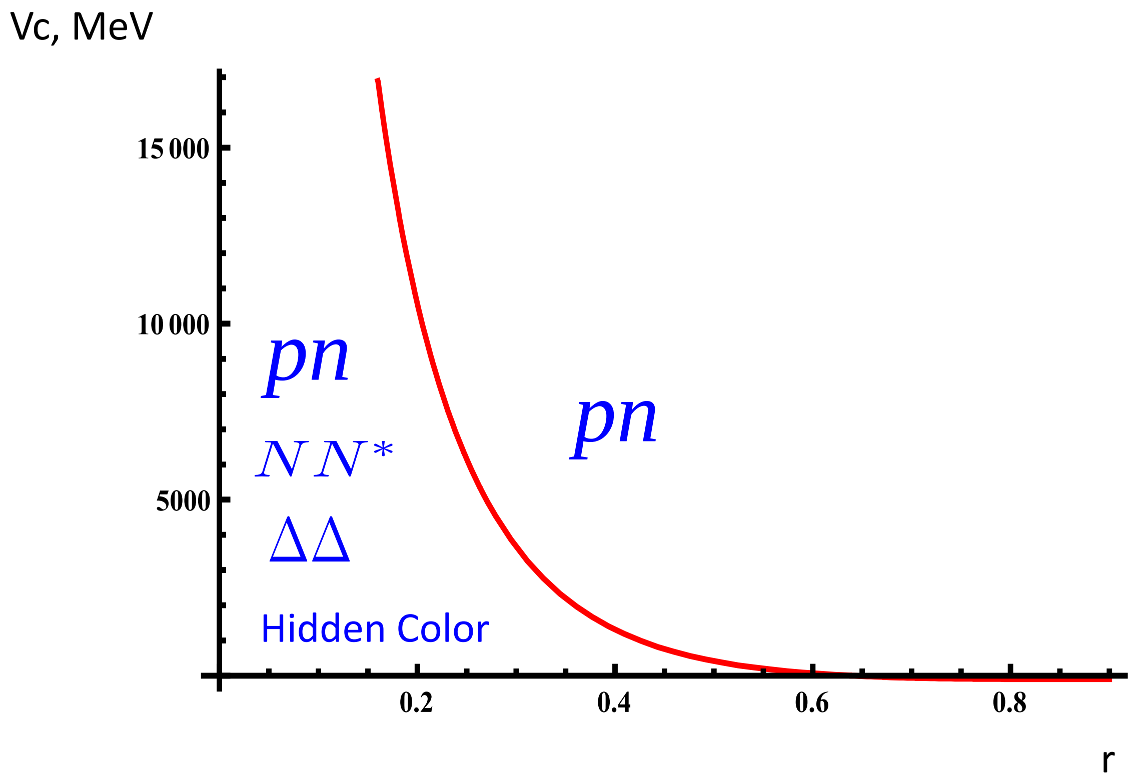
$$\Psi_d = \Psi_{pn} + \Psi_{\Delta\Delta} + \Psi_{NN^*} + \Psi_{hc} \cdots$$

$$\Psi_{hc} = \Psi_{N_c,N_c}$$

$$\Psi^{6q}_{T=0,S=1} = \sqrt{\tfrac{1}{9}}\Psi_{NN} + \sqrt{\tfrac{4}{45}}\Psi_{\Delta\Delta} + \sqrt{\tfrac{4}{5}}\Psi_{CC}$$

The NN core can be due to the orthogonality of

$$\langle \Psi_{N_c,N_c} \mid \Psi_{N,N} \rangle = 0$$



Vc, MeV
15 000
10 000
5000
pn
N N*
ΔΔ
Hidden Color
pn
0.2
0.4
0.6
0.8
r

Nuclear Forces and Nuclear Structure "Standard Approach"

A-body Schroedinger equation interacting through NN -potential

$$\left[-\sum_i \frac{\nabla_i^2}{2m} + \frac{1}{2}\sum_{ij} V(x_i - x_j) + \sum_{ijk} V(x_i, x_j, x_k)\cdots\right]\psi(x_1,\cdots,x_A) = E\psi(x_1,\cdots,x_A)$$

(a)Hartree-Fock potential will smear out main properties NN potential 1990s

$$\left[-\frac{\nabla_N^2}{2m} - V_{HF}(x)\right]\psi_N(x) = E_N\psi_N(x)$$

(b)Ab-Initio Calculations - Modern

$$\left[-\sum_i \frac{\nabla_i^2}{2m} + \frac{1}{2}\sum_{ij} V(x_i - x_j) + \sum_{ijk} V(x_i, x_j, x_k)\cdots\right]\psi(x_1,\cdots,x_A) = E\psi(x_1,\cdots,x_A)$$

Conceptually: How to probe nuclei at short nucleon separations

- Quantum Mechanics allows two or three nucleons to be in short proximity at very short time intervals

- Probe bound nucleons at large internal momenta $p \sim M_N$

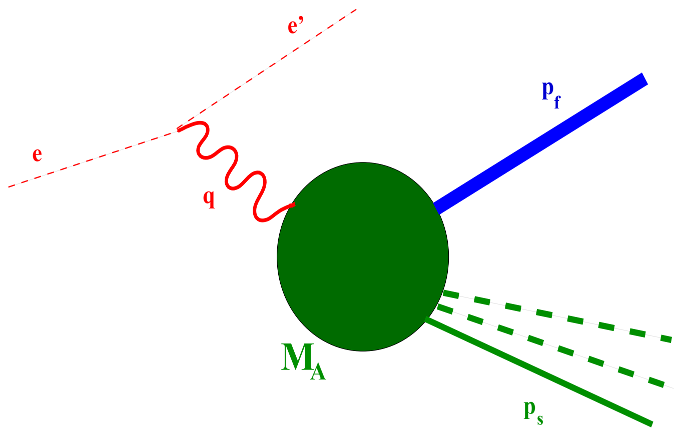
- Due to short range nature of Nuclear Forces nucleons with large internal momenta are in short range NN or NNN correlations (SRCs)

- Need high energy probes to resolve such nucleons in nuclei

Theory of High Energy eA Scattering:

- High-Energy approximations
- Relativism of bound nucleon
- Light-Front Wave Function of Nucleus
- From Schroedinger Equation to Feynman Diagrams
- Emergence of Effective Theory

High Energy Approximations:



$$|\vec{q}| = q_3 \sim p_{f3} \gg p \sim M_N$$

$$Q^2 \geq few \quad \mathrm{GeV}^2$$

Both for QE/DIS

- Emergence of the small parameter

$$\frac{q_-}{q_+} = \frac{q_0 - q_3}{q_0 + q_3} \ll 1 \qquad \mathcal{O}\left(\frac{q_-}{q_+}\right)$$

$$\frac{p_{f-}}{p_{f+}} = \frac{E_f - p_{f3}}{E_f + p_{f3}} \ll 1 \qquad \mathcal{O}\left(\frac{p_{f-}}{p_{f+}}\right)$$

Light-Front Wave Function of the Nucleus

- Emergence of the light- front dynamics

$$\tau = t - z \sim \frac{1}{q_+} \to 0$$

- non relativistic case: due to Galilean relativity observer X can probe all n-nucleons at the same time

$$\Psi(z_1, z_2, z_3, \cdots, z_n, t)$$

- relativistic case: observer X probes all n-nucleons at different n times

$$\Psi(z_1, t_1; z_2, t_2; z_3, t_3 \cdots; z_n, t_n)$$

- observer riding the light-front X probes all n-nucleons at same light-cone time:

$$\Psi_{LF}(\mathcal{Z}_1, \mathcal{Z}_2, \mathcal{Z}_3, \cdots, \mathcal{Z}_n, \tau)$$

$$\tau = t_1 - z_1 = t_2 - z_2 = \cdots = t_n - z_n \qquad \mathcal{Z}_i = t_i + z_i$$

From Schroedinger Equation -> Feynman Diagrams-> Light-Front Wave Function

Schroedinger eq. → **Lipmann-Schwinger Eq.**

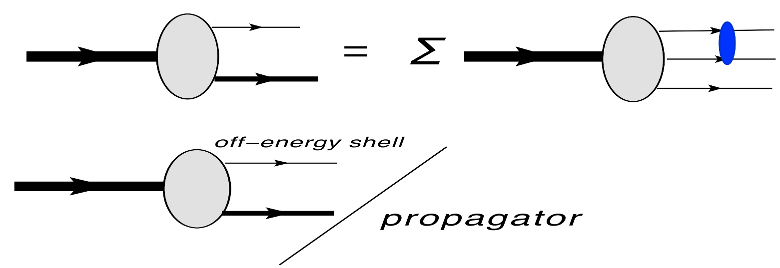
$$\left[-\sum_i \frac{\nabla_i^2}{2m} + \frac{1}{2}\sum_{i,j} V(x_i - x_j)\right]\psi(x_1,\cdots,x_A) = E\psi(x_1,\cdots,x_A)$$

$$\left(\sum_i \frac{k_i^2}{2m} - E_b\right)\Phi(k_1,\cdots,k_A) = -\frac{1}{2}\sum_{i,j}\int U(q)\,\Phi(k_1,\cdots,k_i-q,\cdots k_j+q,\cdots,k_A)d^3q$$

Lipmann-Schwinger Eq → **t- ordered diagrammatic method**

$$\left(\sum_i \frac{k_i^2}{2m} - E_b\right)\Phi(k_1,\cdots,k_A) = -\frac{1}{2}\sum_{i,j}\int U(q)\,\Phi(k_1,\cdots,k_i-q,\cdots k_j+q,\cdots,k_A)d^3q$$

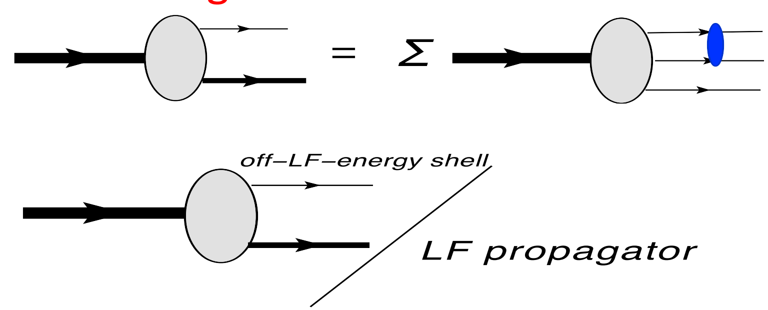
$$\Phi(k_1,\cdots k_A) = \frac{1}{\sum \frac{k_i^2}{2m} - E_b}\Gamma_{A\to N,A-1}$$



Weinberg Eq → **τ - ordered diagrammatic method**

$$\left(\sum \frac{k_{i\perp}^2+m^2}{\alpha_i} - M_A^2\right)\Phi_{LF}(k_1,\cdots k_A) = \frac{1}{2}\sum_{i,j}\int U_{LF}(q)\Phi_{LF}(k_1,\cdots k_A)\prod\frac{d\alpha_i}{\alpha_i}d^2k_{i\perp}$$

$$\Phi_{LF}(k_1,\cdots k_A) = \frac{1}{\sum \frac{k_{i\perp}^2+m^2}{\alpha_i} - M_A^2}\Gamma_{A\to N,A-1}$$



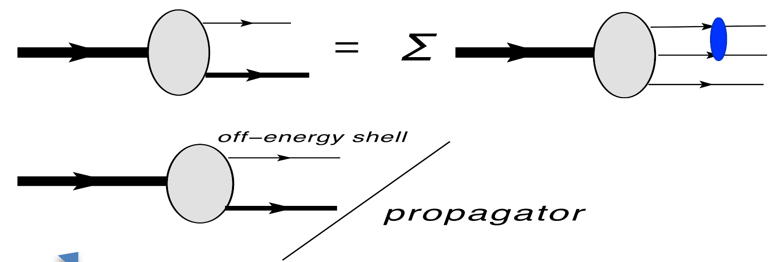
Problem with non-relativistic description: Relativistic Invariance

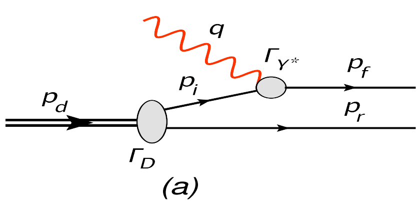
Lipmann-Schwinger Eq →

$$\left(\sum_i \frac{k_i^2}{2m} - E_b\right)\Phi(k_1,\cdots,k_A) = -\frac{1}{2}\sum_{i,j}\int U(q)\,\Phi(k_1,\cdots,k_i-q,\cdots k_j+q,\cdots,k_A)d^3q$$

$$\Phi(k_1,\cdots k_A) = \frac{1}{\sum \frac{k_i^2}{2m} - E_b}\Gamma_{A\to N,A-1}$$

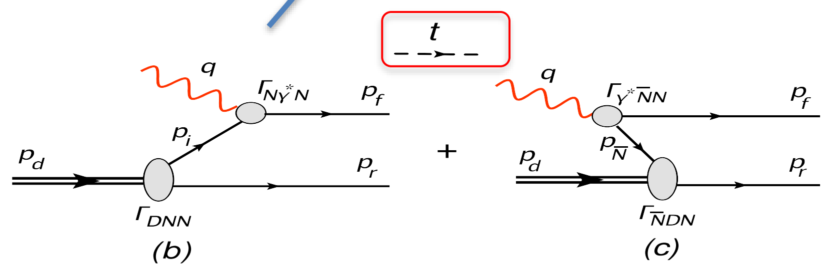
t- ordered diagrammatic method





Feynman/Covariant Diagram

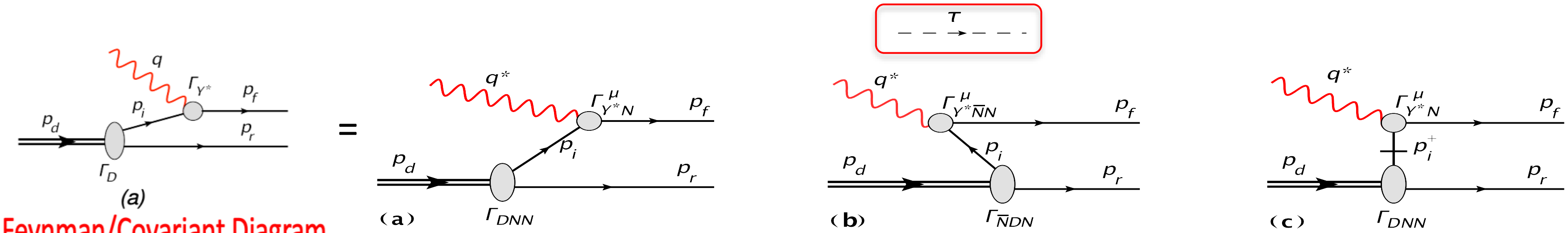
=



$(b) \sim (c)$

Vacuum Fluctuations

Light Front Description: Relativistic Invariance



Feynman/Covariant Diagram

$$(\mathrm{b}) = 0$$

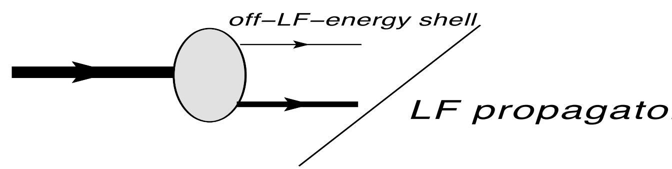
Frank Vera, M.S. PRC 2018

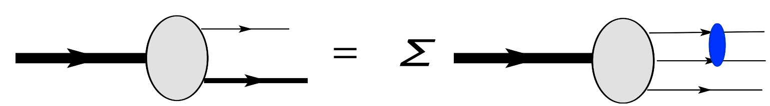
Weinberg Eq →

$$\Phi_{LF}(k_1, \cdots k_A) = \frac{1}{\sum \frac{k_{i\perp}^2 + m^2}{\alpha_i} - M_A^2} \Gamma_{A \to N, A-1}$$

$$\left(\sum \frac{k_{i\perp}^2 + m^2}{\alpha_i} - M_A^2\right)\Phi_{LF}(k_1, \cdots k_A) = \frac{1}{2}\sum_{i,j}\int U_{LF}(q)\Phi_{LF}(k_1, \cdots k_A)\prod \frac{d\alpha_i}{\alpha_i} d^2 k_{i\perp}$$

τ - ordered diagrammatic method





Light-Front Description of the nucleus: Relativistic Invariance

- in the momentum space

$$\Phi_{LF}(k_1,\cdots k_A) = \frac{1}{\sum \frac{k_{i\perp}^2+m^2}{\alpha_i} - M_A^2}\Gamma_{A\to N,A-1}$$

Fourier tranform

$$\Psi_{LF}(\mathcal{Z}_1, \mathcal{Z}_2, \mathcal{Z}_3, \cdots, \mathcal{Z}_n, \tau)$$

$$\Psi_{LF}(\alpha_1, p_{1\perp}; \alpha_2, p_{2\perp}; \alpha_3, p_{3\perp}; \cdots, \alpha_n, p_{n\perp}) \quad \alpha_i = \frac{p_{i-}}{p_{A-}/A}$$

Lorentz boost Invariant

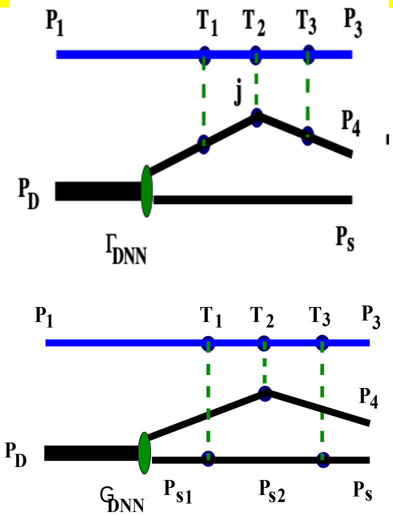
Variables

Proper kinematic variables are not 3d momentum

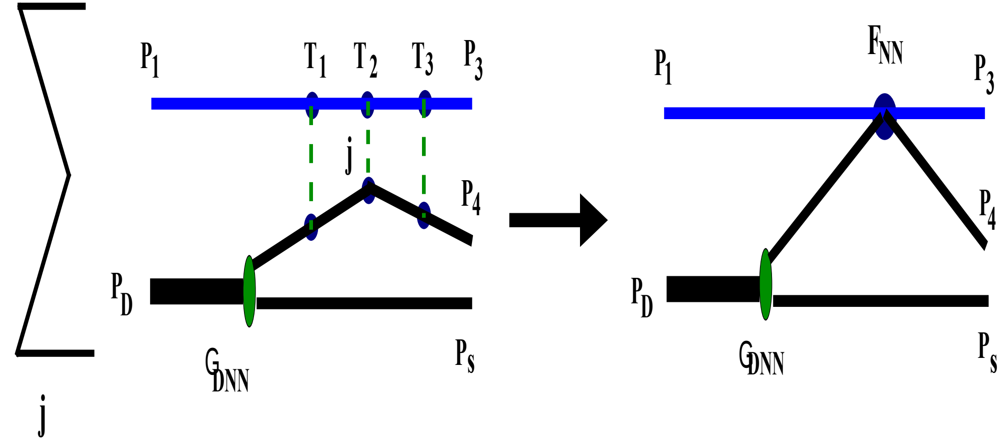
but the Light Front Momentum Fraction: $\alpha = \frac{p_N^+}{p_{NN}^+}$

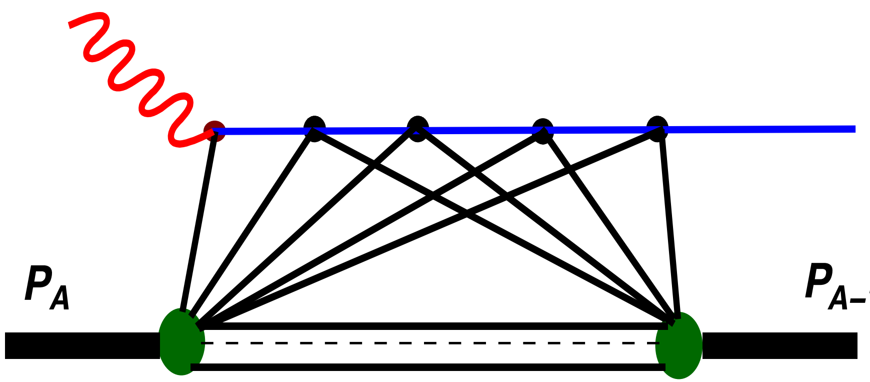
and transverse momentum: $p_\perp$

Emergence of "effective" theory



$\mathcal{O}\left(\frac{p_{1-}}{p_{1+}}\right)$





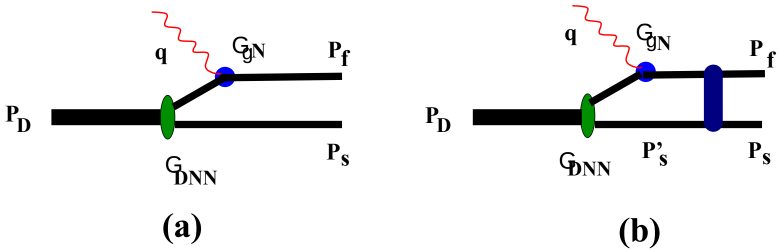
Effective Diagrammatic Rules

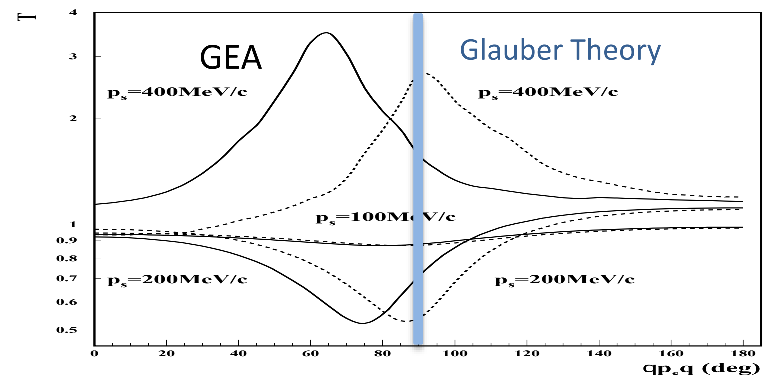
M.S. IJMS 2001

Wave function? $\Phi(k_1, \cdots k_A) = \frac{1}{\sum \frac{k_i^2}{2m} - E_b} \Gamma_{A \to N, A-1}$

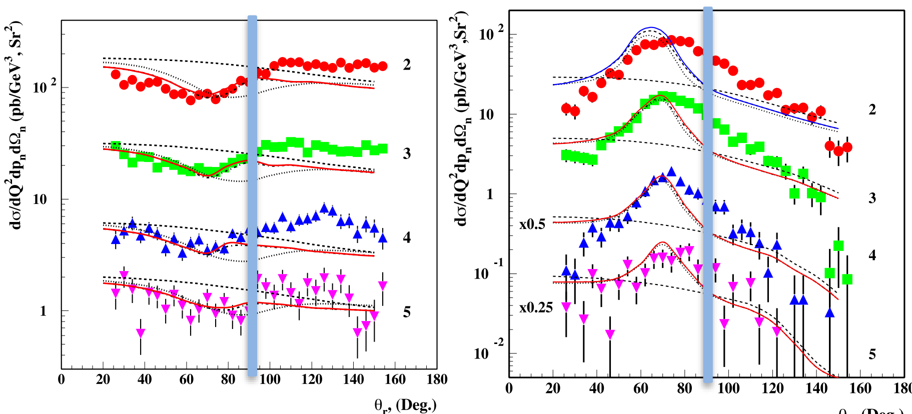
Some Results: $e + d \to e' + p + n$

Frankfurt, M.S., Strikman, PRC 1997

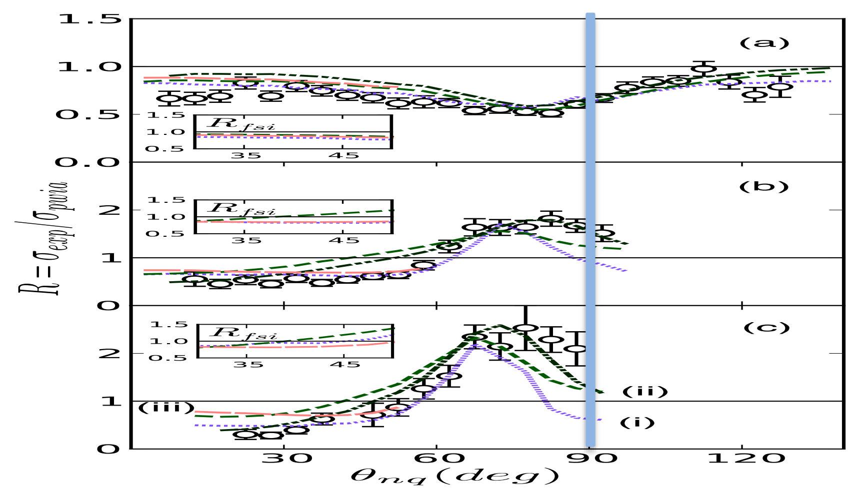




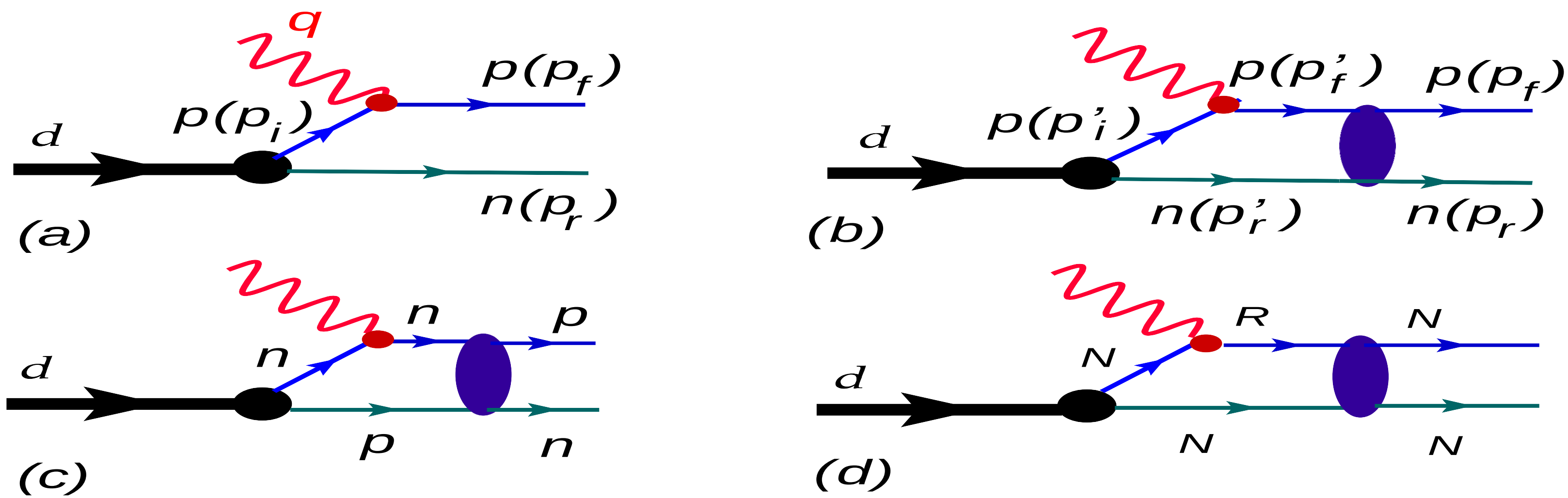
K. Egiyan et al PRL 2008



W. Boeglin et al PRL 2011



M.S. PRC 2010



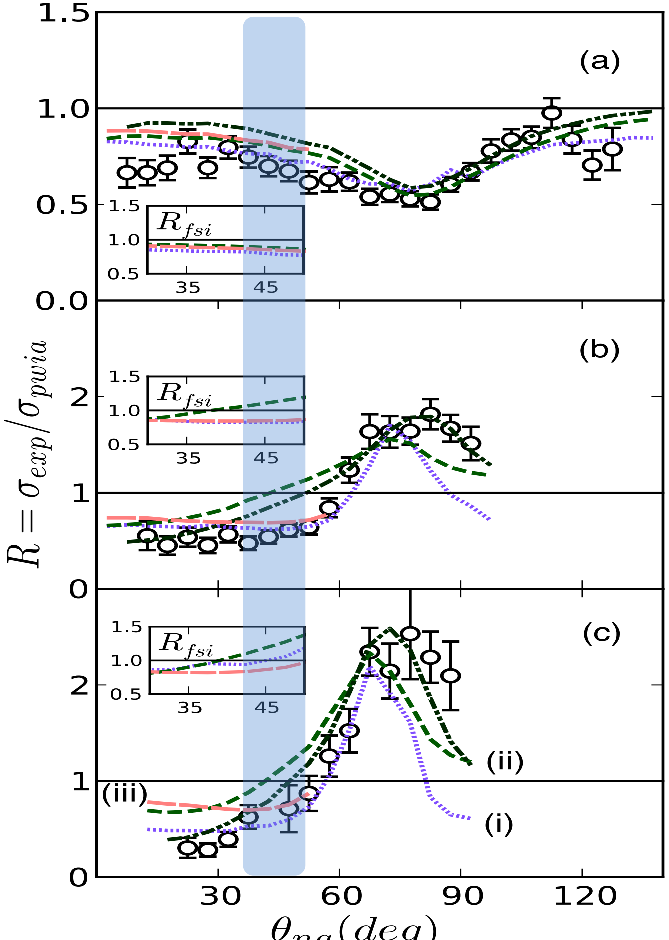
$$\langle s_f, s_r \mid A_0^\mu \mid s_d \rangle = -\bar{u}(p_r, s_r)\Gamma^\mu_{\gamma^* p} \frac{\not{p}_i + m}{p_i^2 - m^2} \cdot \bar{u}(p_f, s_f)\Gamma_{DNN} \cdot \chi^{s_d}$$

$$\langle s_f, s_r \mid A_1^\mu \mid s_d \rangle = -\int \frac{d^4 p'_r}{i(2\pi)^4} \frac{\bar{u}(p_f, s_f)\bar{u}(p_r, s_r) F_{NN}[\not{p}'_r + m][\not{p}_d - \not{p}'_r + \not{q} + m}{(p_d - p'_r + q)^2 - m^2 + i\epsilon}$$

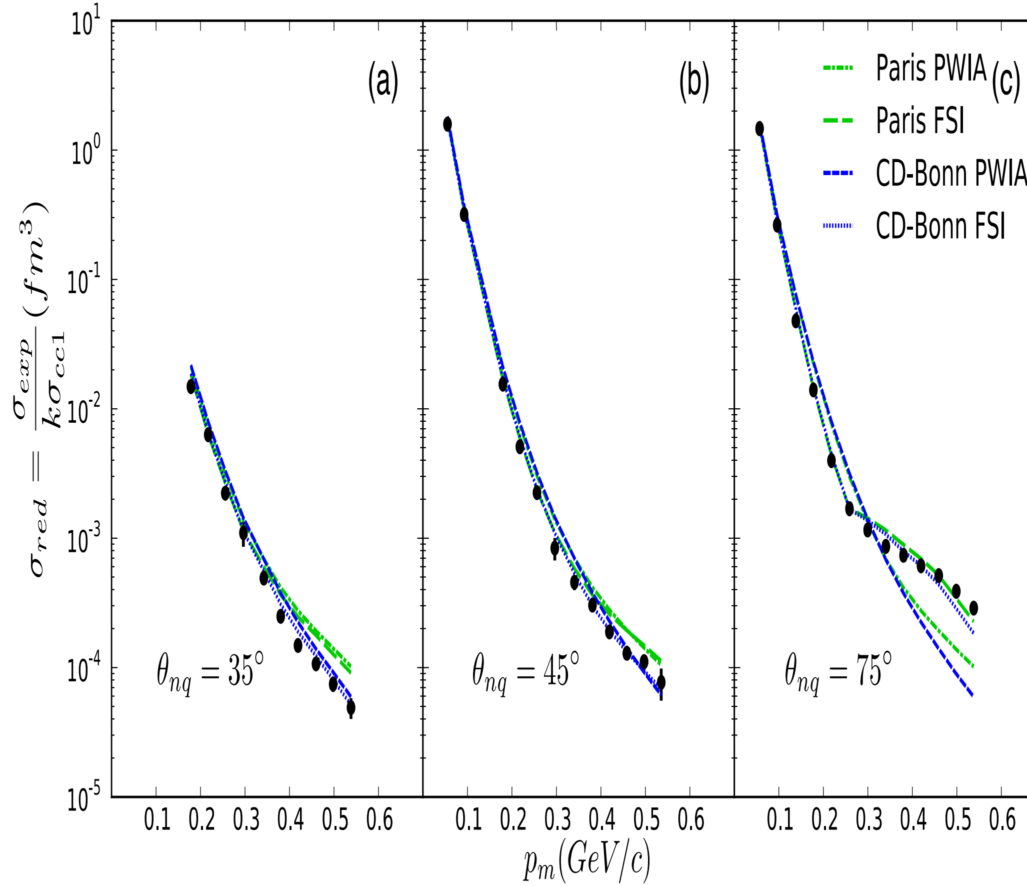
$$\times \frac{\Gamma_{\gamma^* N}[\not{p}_d - \not{p}'_r + m]\Gamma_{DNN}\chi^{s_d}}{((p_d - p'_r)^2 - m^2 + i\epsilon)(p'^2_r - m^2 + i\epsilon)}, \qquad (1$$

M.S. PRC 2010

Probing Deuteron at Small Distances at large Q^2

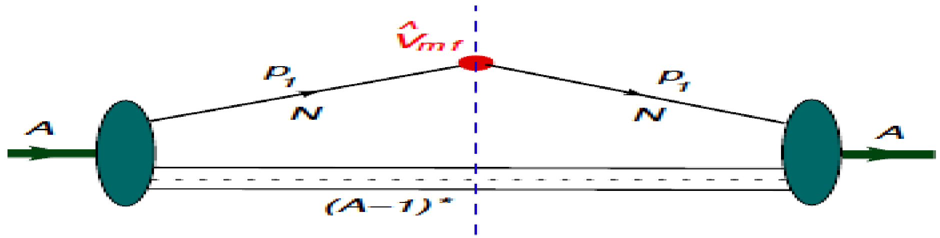


JLab, $Q^2 = 3.5$ GeV^2



Boeglin et al PRL 2011, deuteron probed at up to 550MeV/c

Spectral Function Calculations for Nuclei A>2



$$S_A^{MF} = -Im\int \chi_A^\dagger \Gamma^\dagger_{A,N,A-1}\frac{\not{p}_1+m}{p_1^2-m^2+i\varepsilon}\hat{V}^{MF}\frac{\not{p}_1+m}{p_1^2-m^2+i\times\varepsilon}\left[\frac{G_{A-1}(p_{A-1})}{p_{A-1}^2-M_{A-1}^2+i\varepsilon}\right]^{on}\Gamma_{A,N,A-1}\chi_A\frac{d^4p_{A-1}}{i(2\pi)^4}$$

$$\hat{V}^{MF} = ia^\dagger(p_1,s_1)\delta^3(p_1+p_{A-1})\delta(E_m-E_\alpha)a(p_1,s_1)$$

$$\psi_{N/A}(p_1,s_1,s_A,E_\alpha) = \frac{\bar{u}(p_1,s_1)\Psi^\dagger_{A-1}(p_{A-1},s_{A-1},E_\alpha)\Gamma_{A,N,A-1}\chi_A}{(M_{A-1}^2-p_{A-1}^2)\sqrt{(2\pi)^3 2E_{A-1}}}$$

$$S_A^{MF}(p_1,E_m) = \sum_\alpha\sum_{s_1,s_{A-1}} \mid \psi_{N/A}(p_1,s_1,s_A,E_\alpha)\mid^2 \delta(E_m-E_\alpha)$$

Asymptotics of high momentum component of nuclear Ψ_A

$$\phi_A(k_1,\cdots k_n,\cdots k_A) = \frac{-\frac{1}{2}\int\sum_{i\neq j} V_{ij}(q)\phi_A(k_1,\cdots,k_i-q,\cdots,k_j+q,\cdots,k_A)\frac{d^3q}{(2\pi)^3}}{\sum_{i=1}^{A}\frac{k_i^2}{2m_N}-E_B}$$

$k=p,\ \frac{p^2}{2m}\gg E_B$, $k_i-q=p-q\approx 0\rightarrow q\approx p\quad k_j+q\approx k_j+p\approx 0\rightarrow \boxed{k_i\approx -k_j\approx p}$

$$\phi_A^{(1)}(k_1,\cdots,k_i=p,\cdots,k_j\approx -p,\cdots,k_A)\sim\frac{V_{NN}(p)}{p^2}f(k_1,\cdots'\cdots'\cdots)$$

$$\phi_A^{(2)}(\cdots p,\cdots)\sim\frac{1}{p^2}\int\frac{V_{NN}(q)V_{NN}(p)}{(p-q)^2}d^3q$$

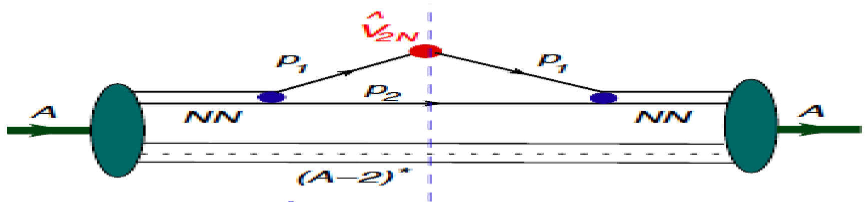
$V_{NN}=q^{-n}$ with $n>1$ $\quad\phi_A^{(2)}(\cdots p,\cdots)\sim\frac{V(p)}{p^2}\int_{q_{min}}^{\infty}\frac{dq}{q^n}$

and is finite range $\quad q_{min}$

$$\phi_A^{(2)}\ll\phi_A^{(1)}$$

Frankfurt, Strikman 1988
Frankfurt, MS, Strikman 2008

2N SRC model



$$P^N_{A,2N}(\alpha_1, p_{1,\perp}, \tilde{M}^2_N) = \sum_{s_2, s_{NN}, s_{A-2}} \int \chi^\dagger_A \Gamma^\dagger_{A\to NN,A-2} \chi_{A-2}(p_{A-2}, s_{A-2})$$
$$\times \frac{\chi_{NN}(p_{NN}, s_{NN})\chi^\dagger_{NN}(p_{NN}, s_{NN})}{p^2_{NN} - M^2_{NN}} \Gamma^\dagger_{NN\to NN} \frac{u(p_1, s_1)u(p_2, s_2)}{p^2_1 - M^2_N}$$
$$\times \left[2\alpha^2_1 \delta(\alpha_1 + \alpha_2 + \alpha_{A-2} - A)\delta^2(p_{1,\perp} + p_{2,\perp} + p_{A-2,\perp})\delta(\tilde{M}^2_N - \tilde{M}^{(2N),2}_N)\right] \frac{\bar{u}(p_1, s_1)\bar{u}(p_2, s_2)}{p^2_1 - M^2_N}$$
$$\times \Gamma_{NN\to NN} \frac{\chi_{NN}(p_{NN}, s_{NN})\chi^\dagger_{NN}(p_{NN}, s_{NN})}{p^2_{NN} - M^2_{NN}} \chi^\dagger_{A-2}(p_{A-2}, s_{A-2})\Gamma_{A,NN,A-2}\chi_A$$
$$\times \frac{d\alpha_2}{\alpha_2} \frac{d^2 p_{2,\perp}}{2(2\pi)^3} \frac{d\alpha_{A-2}}{\alpha_{A-2}} \frac{d^2 p_{A-2,\perp}}{2(2\pi)^3}. \quad (1)$$

O. Artiles & M.S. Phys. Rev. C 2016

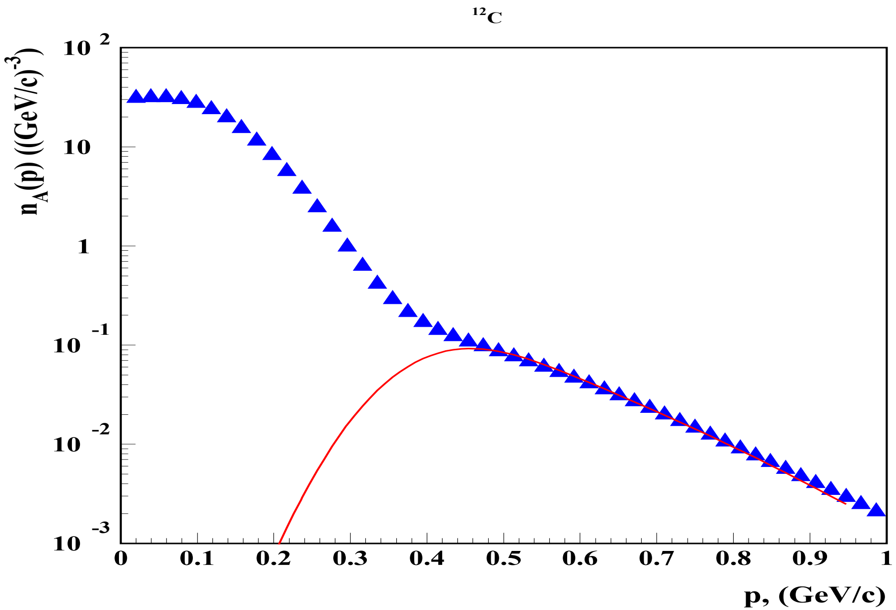
$$\rho_A(\alpha_N, p_{N,\perp}) = \int P_A(\alpha_N, p_{N,\perp}, \tilde{M}^2_N)\tfrac{1}{2} d\tilde{M}^2_N$$

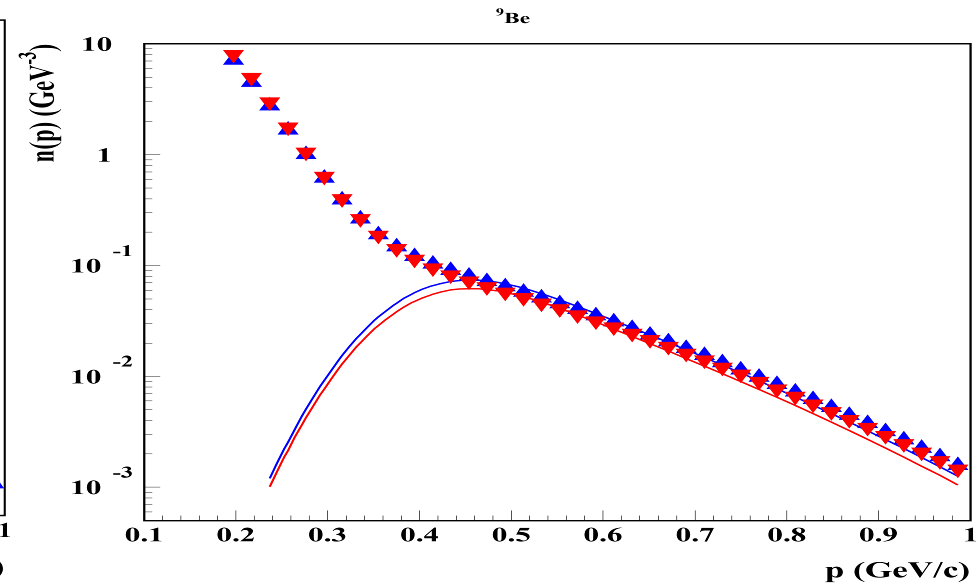
$$\psi^{LF}_{pn}(\alpha, p_\perp) \approx C\psi^{LF}_d(\alpha, p_\perp)$$

$$\psi^{s_{NN}}_{2N}(\beta_1, k_{1,\perp}, s_1, s_2) = -\frac{1}{\sqrt{2(2\pi)^3}} \frac{\bar{u}(p_1,s_1)\bar{u}(p_2,s_2)\Gamma_{NN\to NN}\cdot\chi_{NN}(p_{NN},s_{NN})}{\frac{1}{2}[M^2_{NN} - 4(M^2_N + k^2_1)]}$$

$$\psi_{CM}(\alpha_{NN}, k_{NN,\perp}) = -\frac{1}{\sqrt{\frac{A-2}{2}}} \frac{1}{\sqrt{2(2\pi)^3}} \frac{\chi^\dagger_{NN}(p_{NN},s_{NN})\chi^\dagger_{A-2}(p_{A-2},s_{A-2})\Gamma_{A\to NN,A-2}\chi^{s_A}_A}{\frac{2}{A}[M^2_A - s_{NN,A-2}(k_{CM})]}$$

2N SRC model Non Relativistic Approximation





1. Short-Range NN Correlations

$$\text{for large } k > k_{Fermi}$$

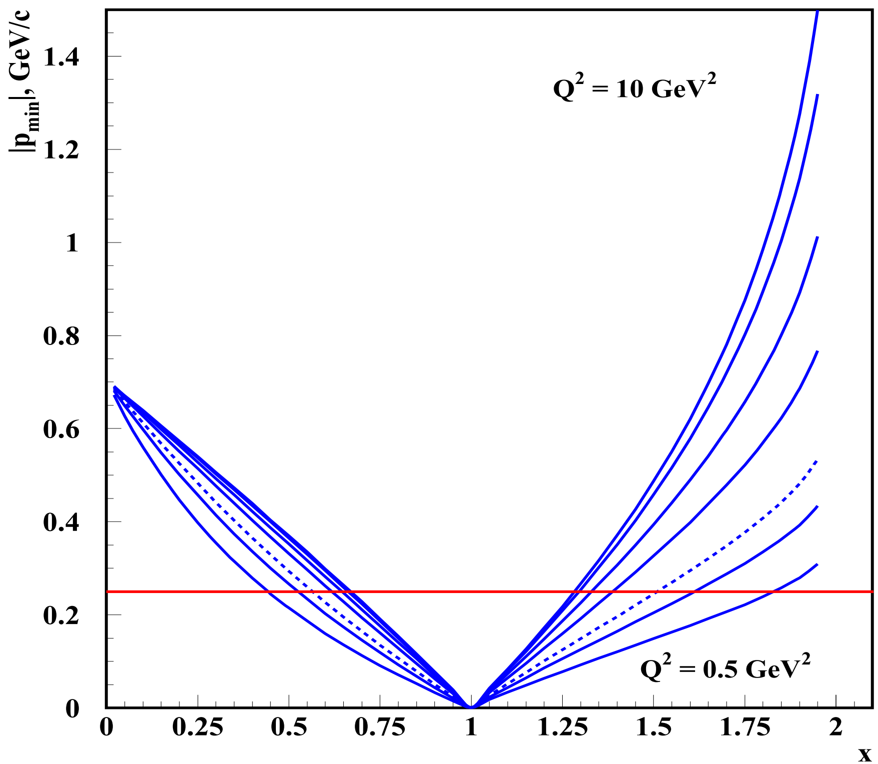
$$n_A(k) \approx a_{NN}(A) n_{NN}(k)$$

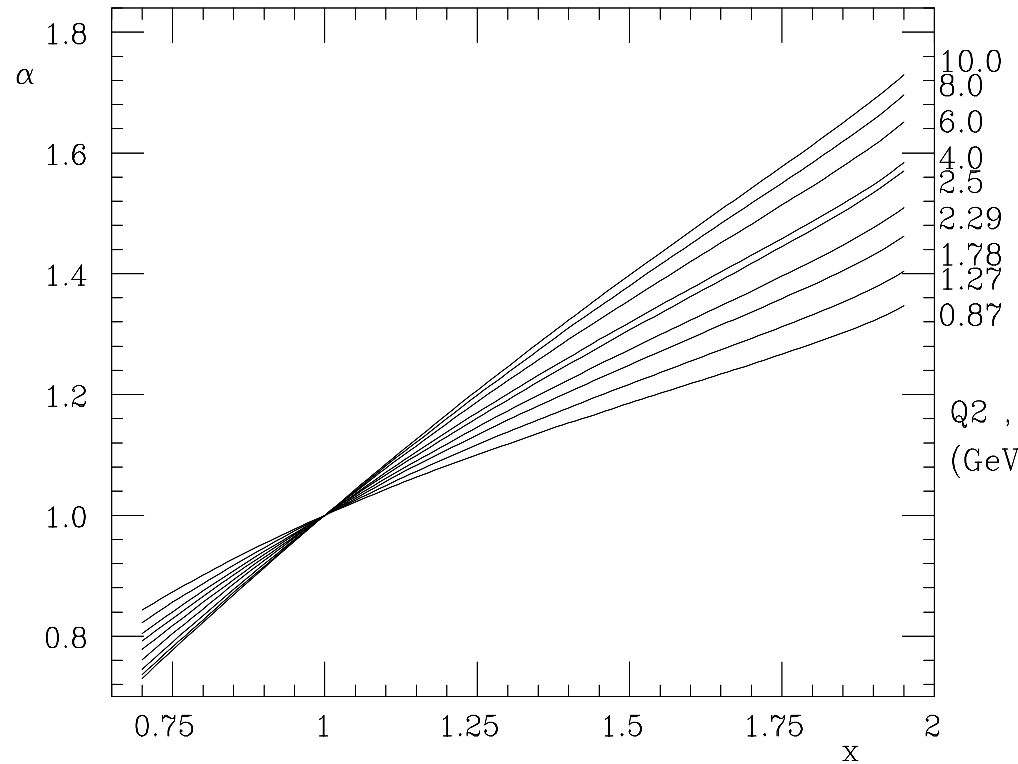
Frankfurt, Strikman Phys. Rep, 1988
Day, Frankfurt, Strikman, MS, Phys. Rev. C 1993

- Experimental observations

Egiyan et al, 2002,2006
Fomin et al, 2011

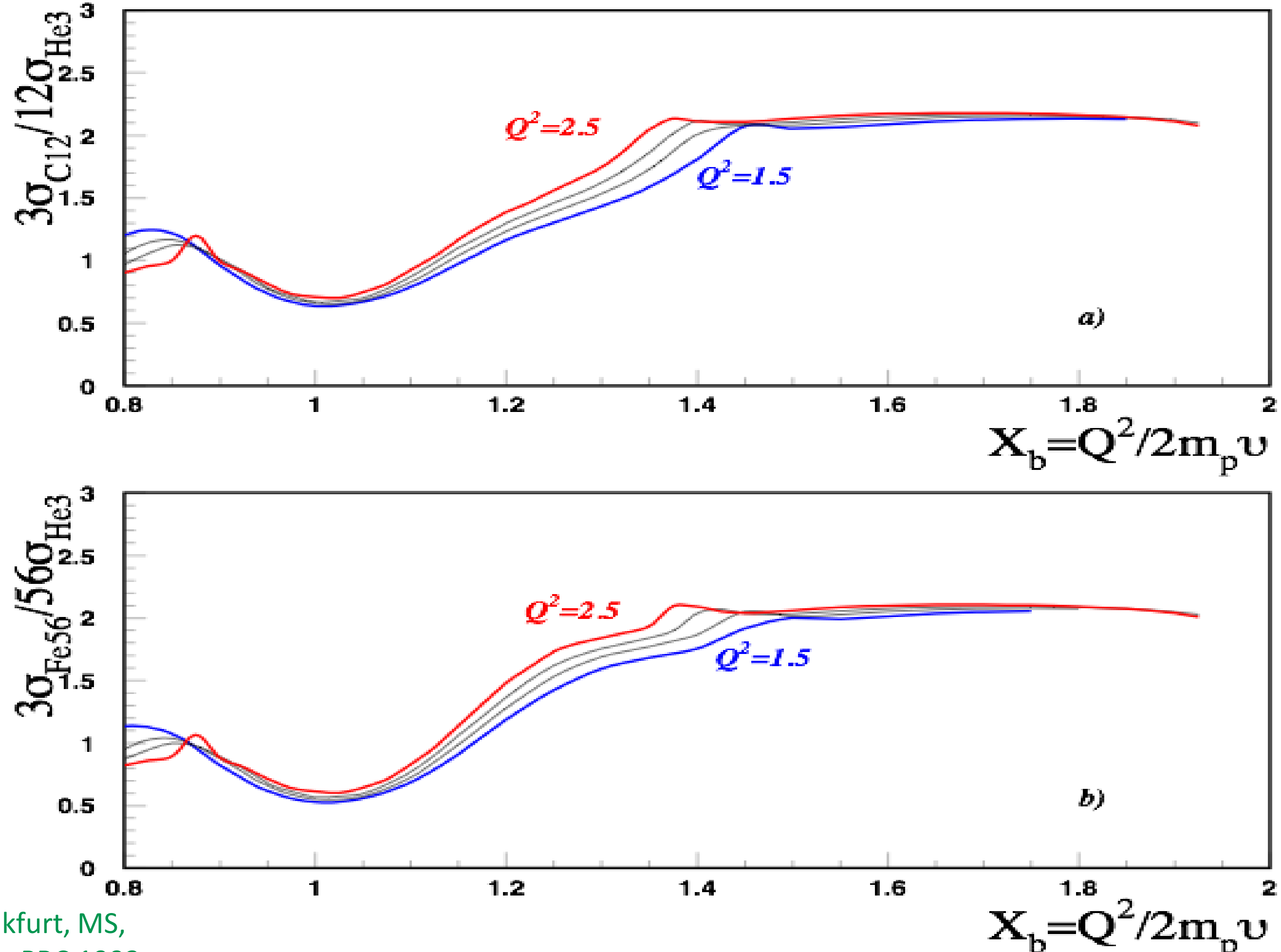
Nuclear Scaling in QE Inclusive A(e,e')X reaction at x>1 region





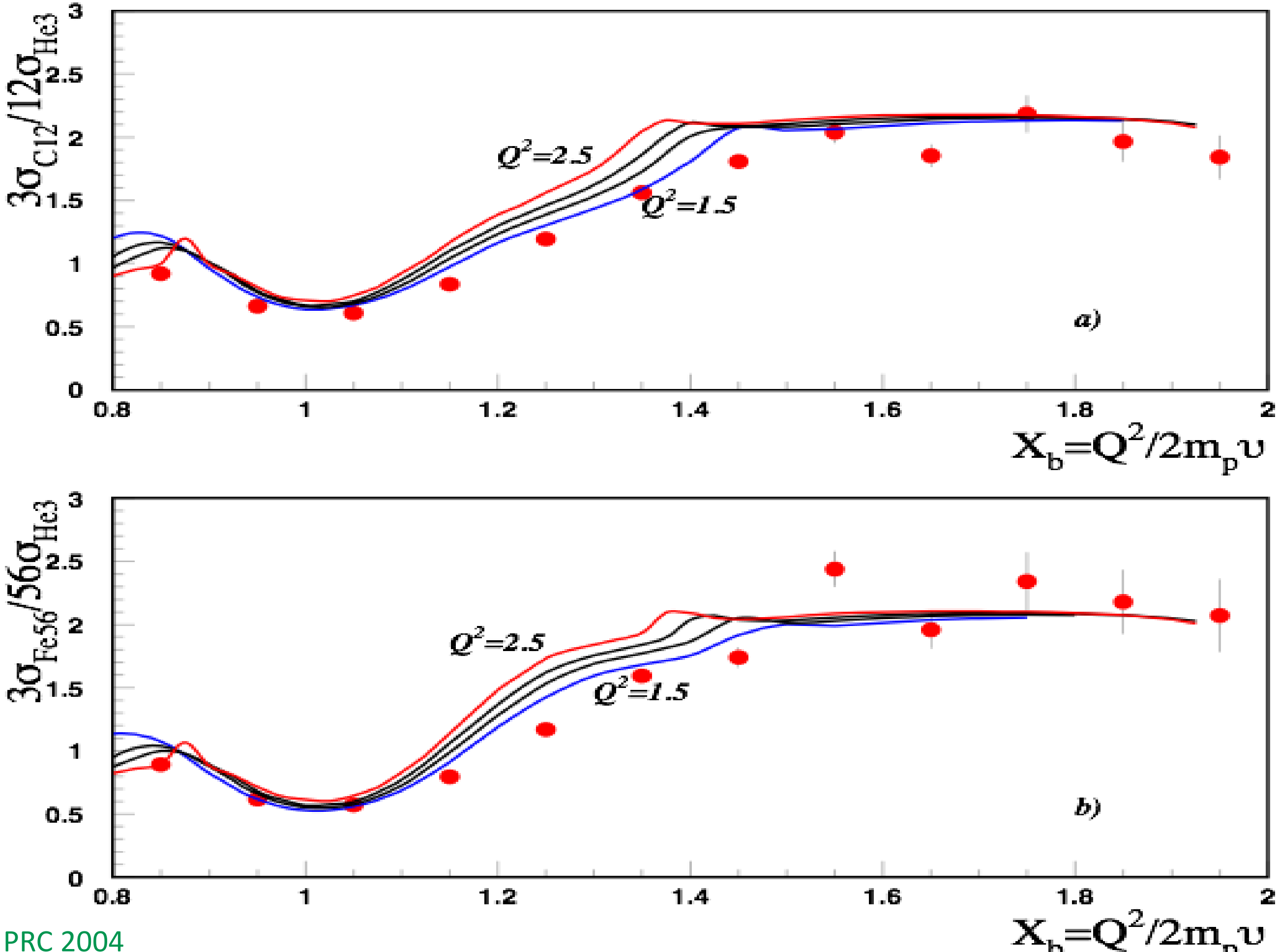
Day, Frankfurt, MS, Strikman,
PRC 1993

A(e,e')

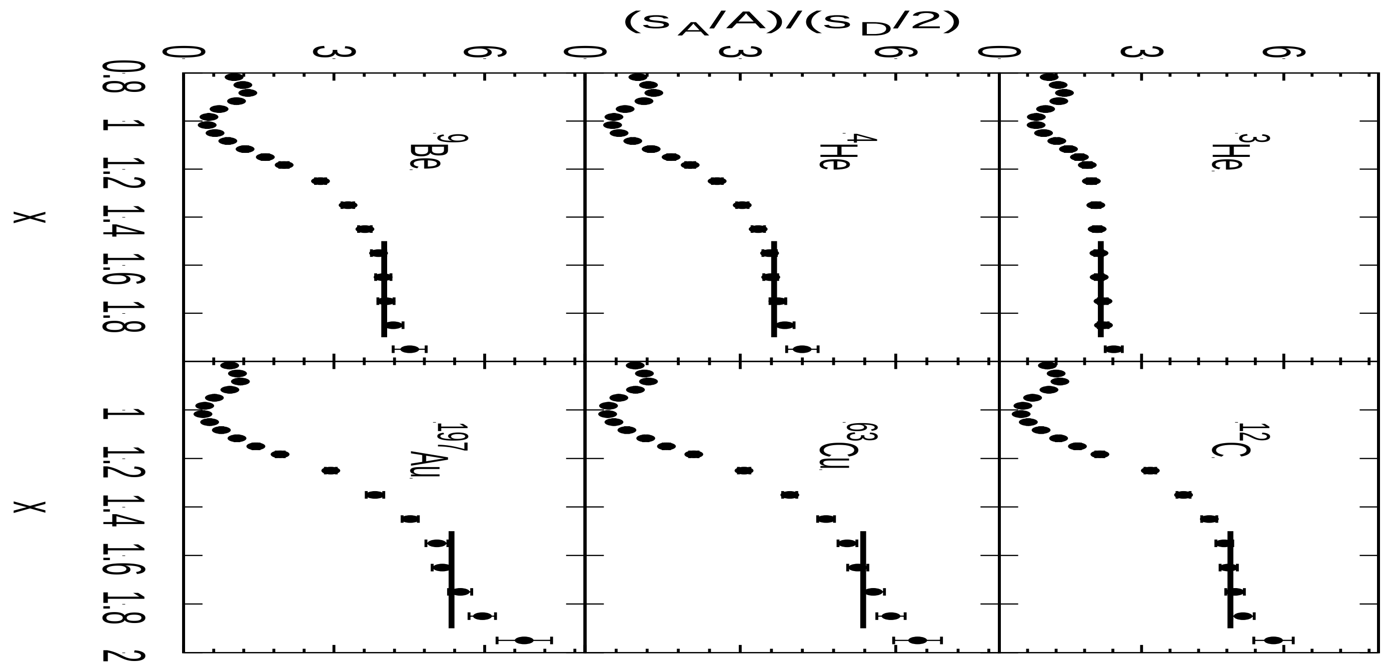


Day, Frankfurt, MS,
Strikman, PRC 1993

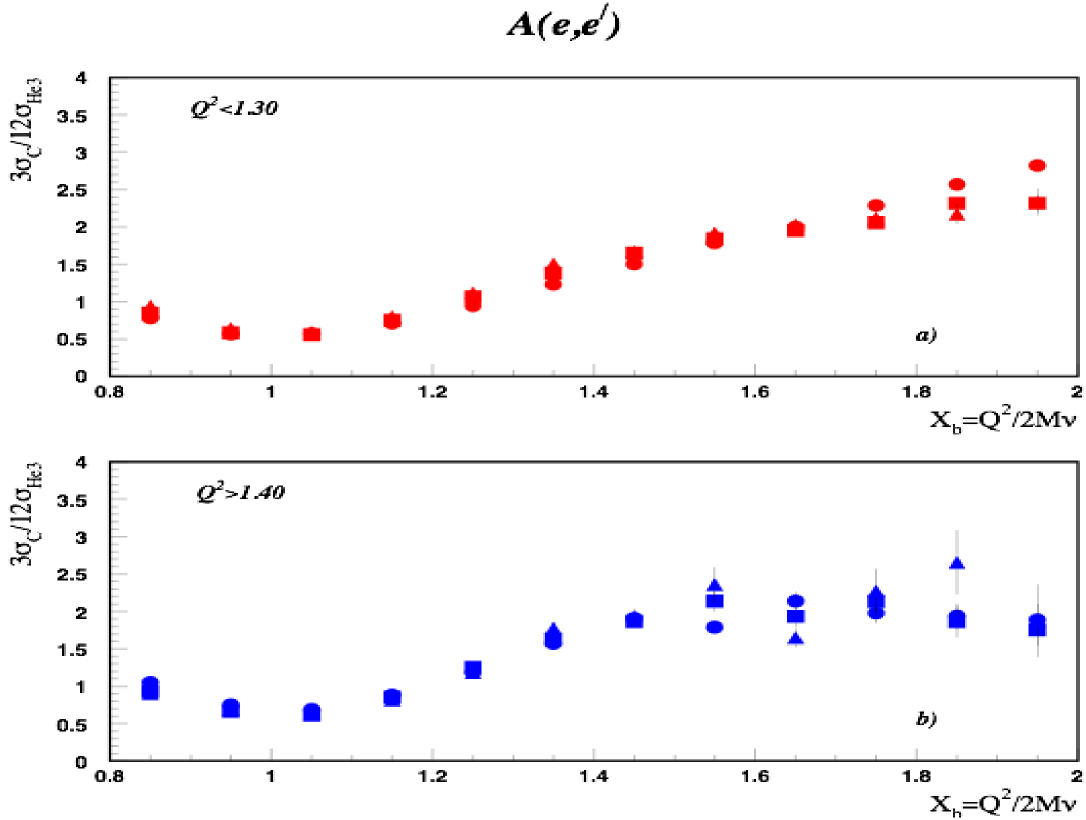
$A(e,e')$



Egiyan, et al PRC 2004



Fomin et al PRL 2011


A(e,e/)
Q²<1.30
a)
3σ_C/12σ_He3
X_b=Q²/2Mν
Q²>1.40
b)
3σ_C/12σ_He3
X_b=Q²/2Mν

2N SRCs:

Proper Variables of 2N SRC are

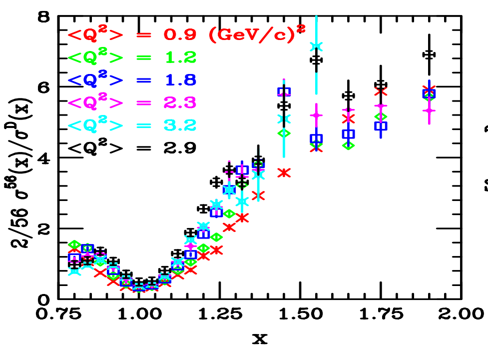
- the Light Front Momentum Fraction: $\alpha = \frac{p_N^+}{p_{NN}^+}$
- transverse momentum: $p_\perp$

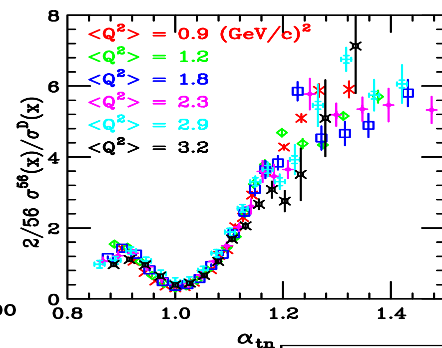
Back to inclusive A(e,e')X scattering

$$\sigma_{eA} = \sum \sigma_{eN} \cdot \rho_A(\alpha) \quad \text{where} \quad \rho_A(\alpha) = \int \rho_A(\alpha, p_\perp) d^2 p_\perp$$

$$1.3 \leq \alpha_{2N} \leq 1.5$$

$$\alpha_{2N} = 2 - \frac{q_- + 2m_N}{2m_N}\left(1 + \frac{\sqrt{W_{2N}^2 - 4m_N^2}}{W_{2N}}\right)$$

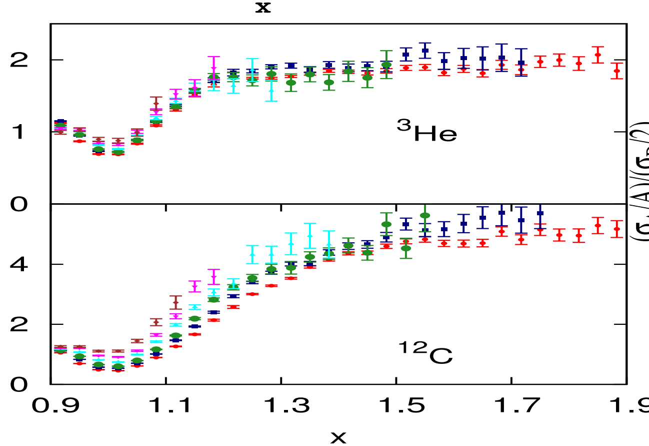


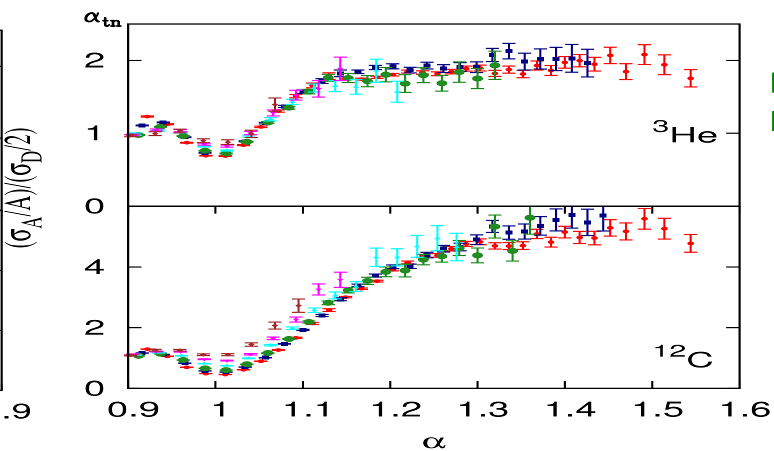


$$\alpha \,\Big|_{Q^2 \to \infty} \to x$$

$$\alpha \,\Big|_{x \to 1} \to 1$$

J.Arrington, D.Higinbotham
G.Rosner, M.S. Prog. PNP 2012





N.Fomin, D.Higinbotham
M.S., P.Sovignon ARNPS, 2017

The Meaning of scaling values

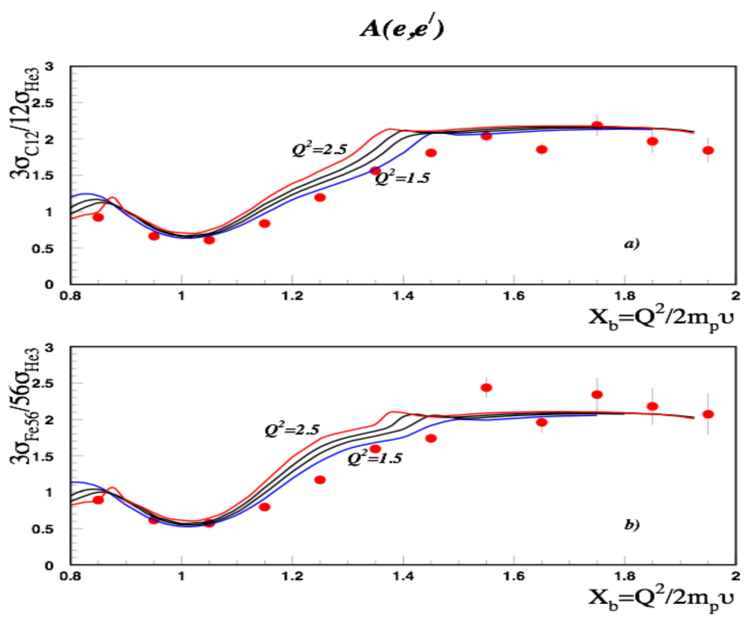
Day, Frankfurt, MS, Strikman, PRC 1993

Frankfurt, MS, Strikman, IJMP A 2008

Fomin et al PRL 2011

$$R = \frac{A_2\sigma[A_1(e,e')X]}{A_1\sigma[A_2(e,e')X]}$$

$$\text{For } 1 < x < 2 \ R \approx \frac{a_2(A_1)}{a_2(A_2)}$$



Egiyan, et al PRL 2006, PRC 2004

a2's as relative probability of 2N SRCs

Table 1: The results for $a_2(A, y)$

A	y	This Work	Frankfurt et al	Egiyan et al	Famin et al
^{3}He	0.33	2.07±0.08	1.7±0.3		2.13±0.04
^{4}He	0	3.51±0.03	3.3±0.5	3.38±0.2	3.60±0.10
^{9}Be	0.11	3.92±0.03			3.91±0.12
^{12}C	0	4.19±0.02	5.0±0.5	4.32±0.4	4.75±0.16
^{27}Al	0.037	4.50±0.12	5.3±0.6		
^{56}Fe	0.071	4.95±0.07	5.6±0.9	4.99±0.5	
^{64}Cu	0.094	5.02±0.04			5.21±0.20
^{197}Au	0.198	4.56±0.03	4.8±0.7		5.16±0.22

2. Dominance of the (pn) component of SRC

for large $k > k_{Fermi}$

$$n_A(k) \approx a_{NN}(A) n_{NN}(k)$$

$$P_{pn/pX} = 0.92^{+0.08}_{-0.18}$$

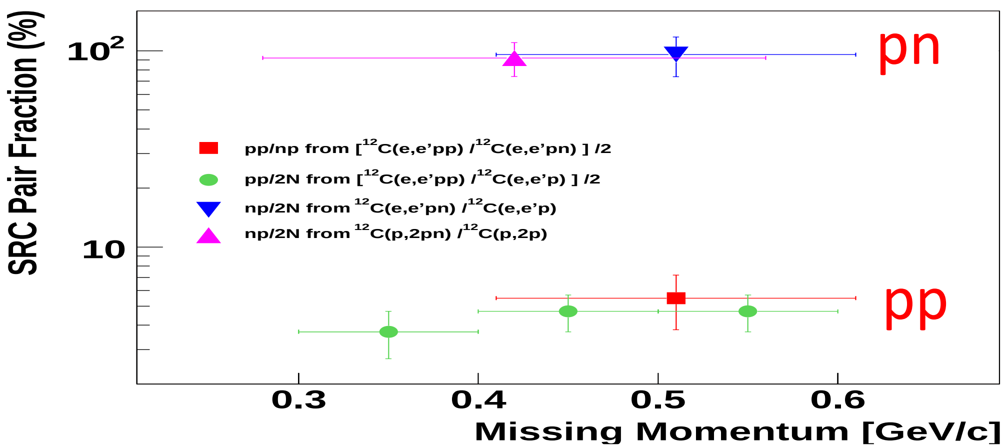
Theoretical analysis of BNL Data $A(p,2p)X$ reaction

$$\frac{P_{pp}}{P_{pn}} \leq \frac{1}{2}(1 - P_{pn/pX}) = 0.04^{+0.09}_{-0.04}.$$

E. Piasetzky, MS, L. Frankfurt, M. Strikman, J.Watson PRL , 2006

$$P_{pp/pn} = 0.056 \pm 0.018$$

Direct Measurement at JLab R.Subdei, et al Science , 2008



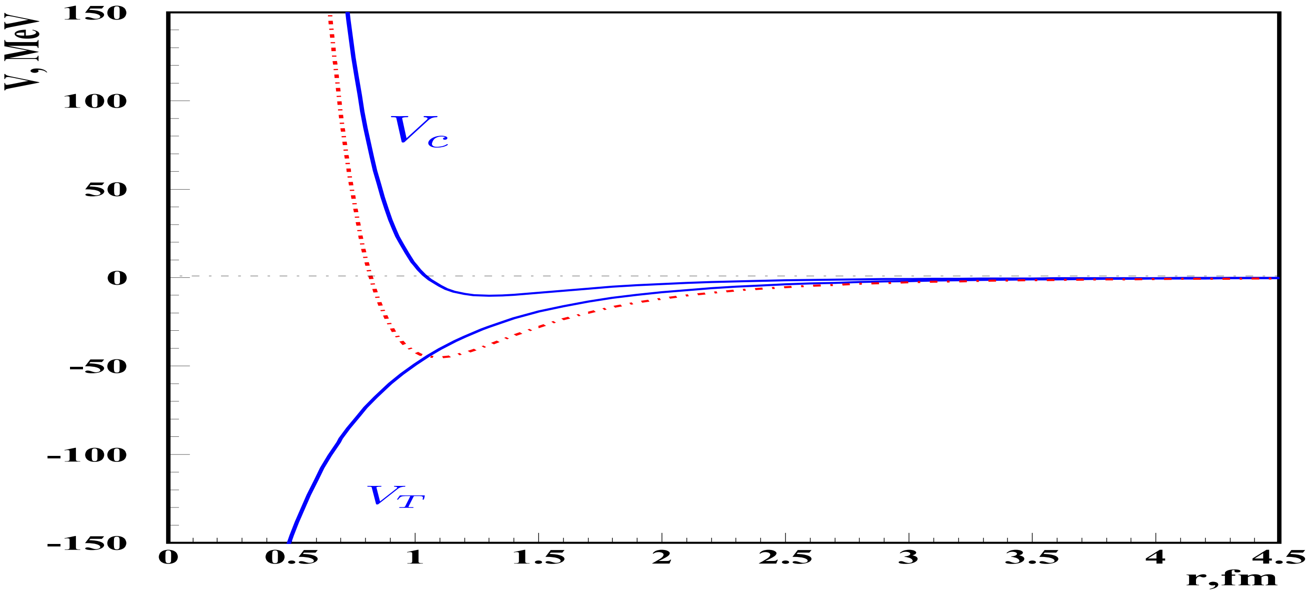
Factor of 20

Expected 4
(Wigner counting)

Theoretical Interpretation

$$\phi_A^{(1)}(k_1,\cdots,k_i=p,\cdots,k_j\approx -p,\cdots,k_A)\sim \frac{V_{NN}(p)}{p^2}f(k_1,\cdots'\cdots'\cdots\quad)$$

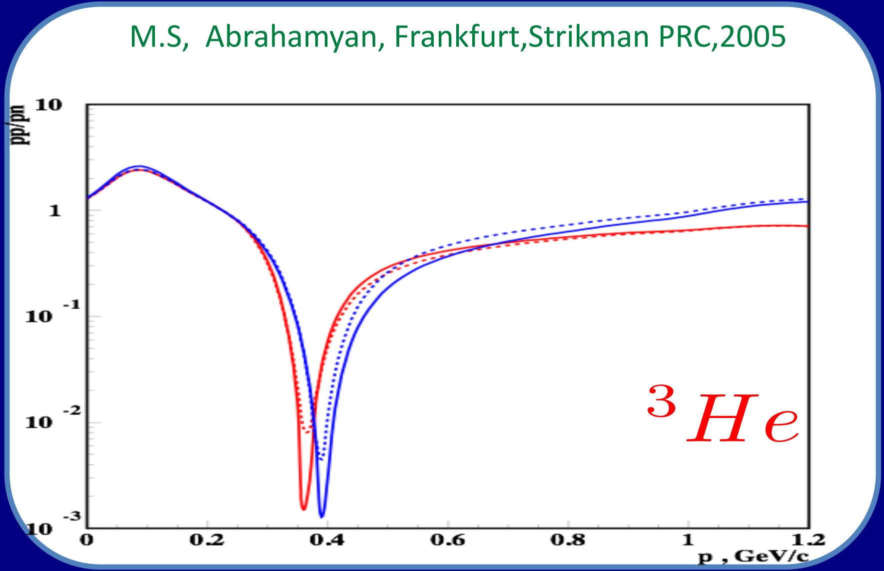
$$n_A(k)\approx a_{NN}(A)n_{NN}(k)$$



Explanation lies in the dominance of the <u>tensor</u> part in the NN interaction

$$V_{NN}(r) \approx V_c(r) + V_t(r) \cdot S_{12}(r) + V_{LS} \cdot \vec{L}\vec{S}$$

$$S_{12} = 3(\sigma_1 \cdot \hat{r})(\sigma_2 \cdot \hat{r}) - \sigma_1\sigma_2$$

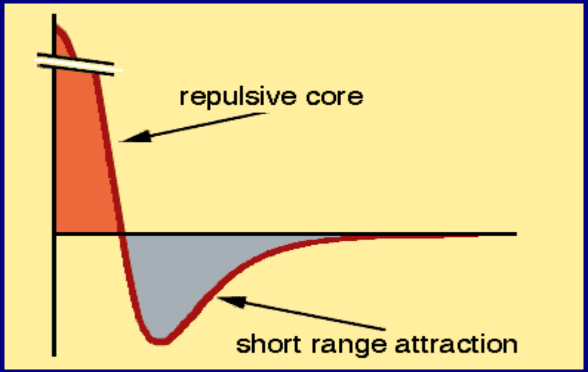


$$S_{12}|pp\rangle = 0$$

$$S_{12}|nn\rangle = 0$$ Isospin 1 states

$$S_{12}|pn\rangle = 0$$

$$S_{12}|pn\rangle \neq 0$$ Isospin 0 states



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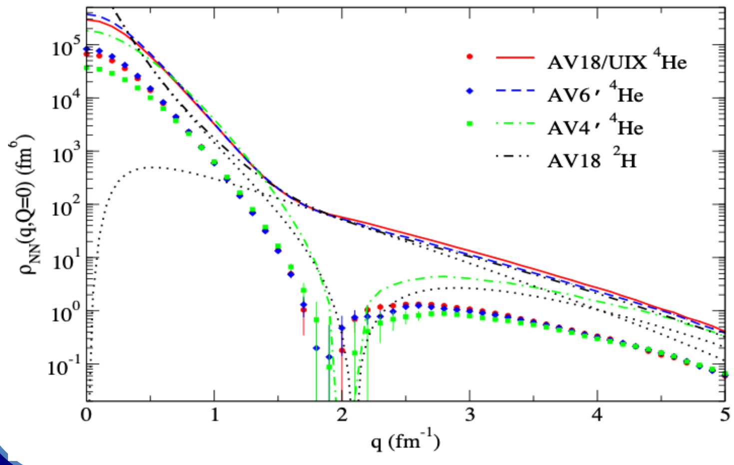
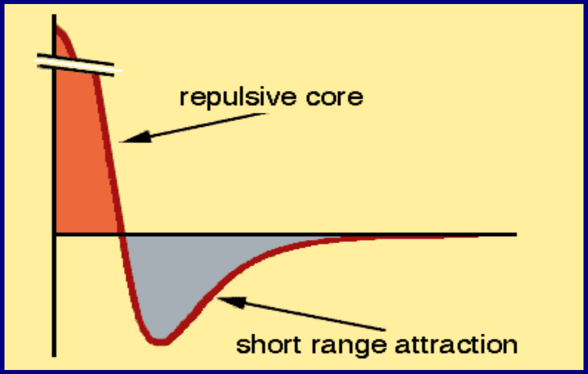
$S_{12}|pn\rangle \neq 0$ Isospin 0 states

Sciavilla, Wiringa, Pieper, Carlson PRL,2007





- *Dominance of pn short range correlations as compared to pp and nn SRCS*

2006-2008s

- *Dominance of NN Tensor as compared to the NN Central Forces at <= 1fm*

2011- present

- **Two New Properties of High Momentum Component**

- **Energetic Protons in Neutron Rich Nuclei**

3. Momentum sharing properties of Nuclear High Momentum Component

at $p > k_F$

$$n^A(p) \sim a_{NN}(A) \cdot n_{NN}(p)$$

- Dominance of pn Correlations
  (neglecting pp and nn SRCs)

$$n_{NN}(p) \approx n_{pn}(p) \approx n_{(d)}(p)$$

$$n^A(p) \sim a_{pn}(A) \cdot n_d(p)$$

$$a_2(A) \equiv a_{NN}(A) \approx a_{pn}(A)$$

MS,arXiv: 1102.3973, 1210.3280,2011-12
Phys. Rev. C 2014

- Define momentum distribution of proton & neutron

$$n^A(p) = \frac{Z}{A} n_p^A(p) + \frac{A-Z}{A} n_n^A(p)$$

$$\int n_{p/n}^A(p) d^3p = 1$$

- Define

$$I_p = \frac{Z}{A} \int_{k_F}^{600} n_p^A(p) d^3p \qquad I_n = \frac{A-Z}{A} \int_{k_F}^{600} n_n^A(p) d^3p$$

- and observe that in the limit of no pp and nn SRCs

$$\boxed{I_p = I_n}$$

- Neglecting CM motion of SRCs

$$\frac{Z}{A} n_p^A(p) \approx \frac{A-Z}{A} n_n^A(p)$$

First Property: Approximate Scaling Relation

-if contributions by pp and nn SRCs are neglected and the pn SRC is assumed at rest

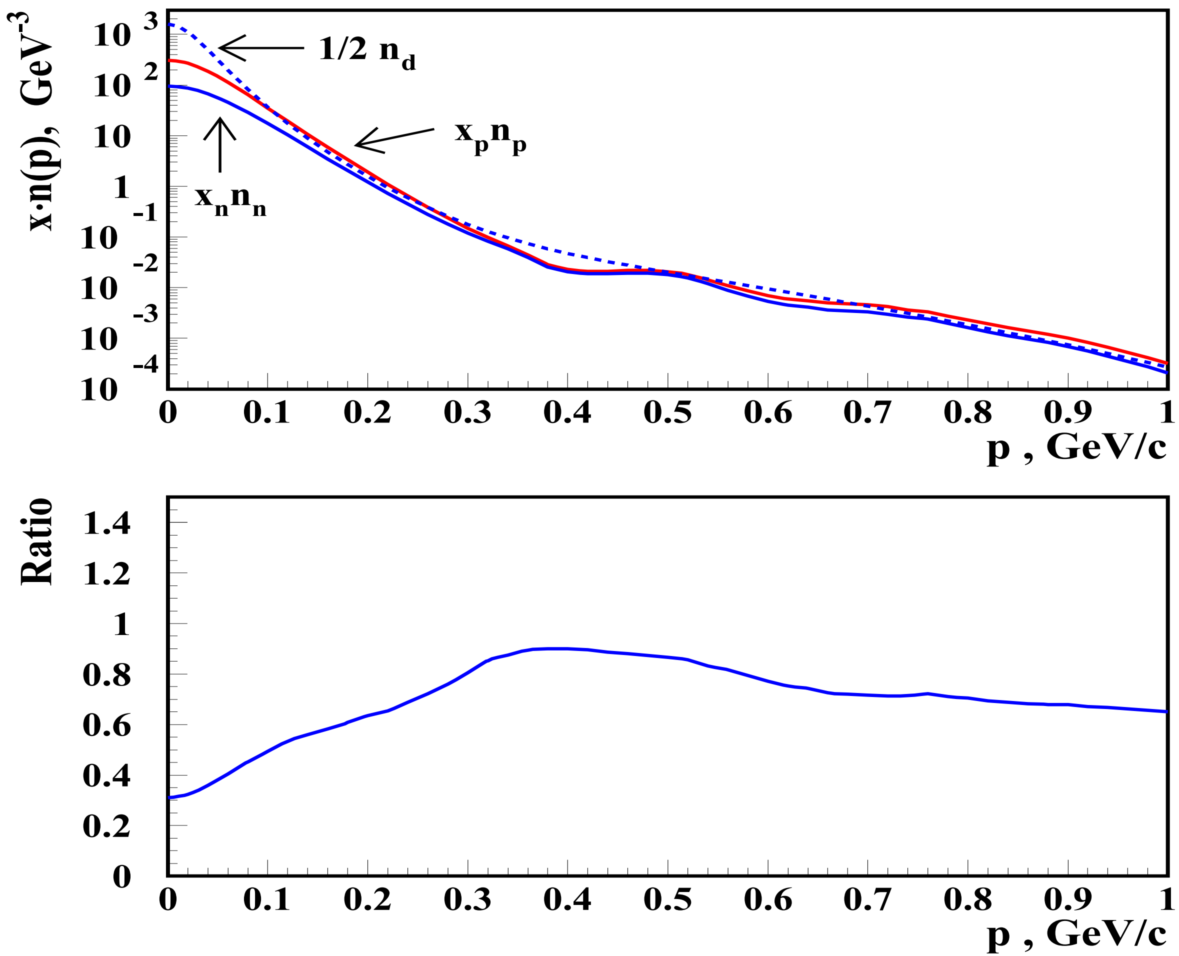
MS,arXiv:1210.3280
Phys. Rev. C 2014

- for $\sim k_F - 600$ MeV/c region:

$$x_p \cdot n_p^A(p) \approx x_n \cdot n_n^A(p) = \frac{a_{pn}}{2}(A,Z) n_d(p)$$

where $x_p = \frac{Z}{A}$ and $x_n = \frac{A-Z}{A}$.

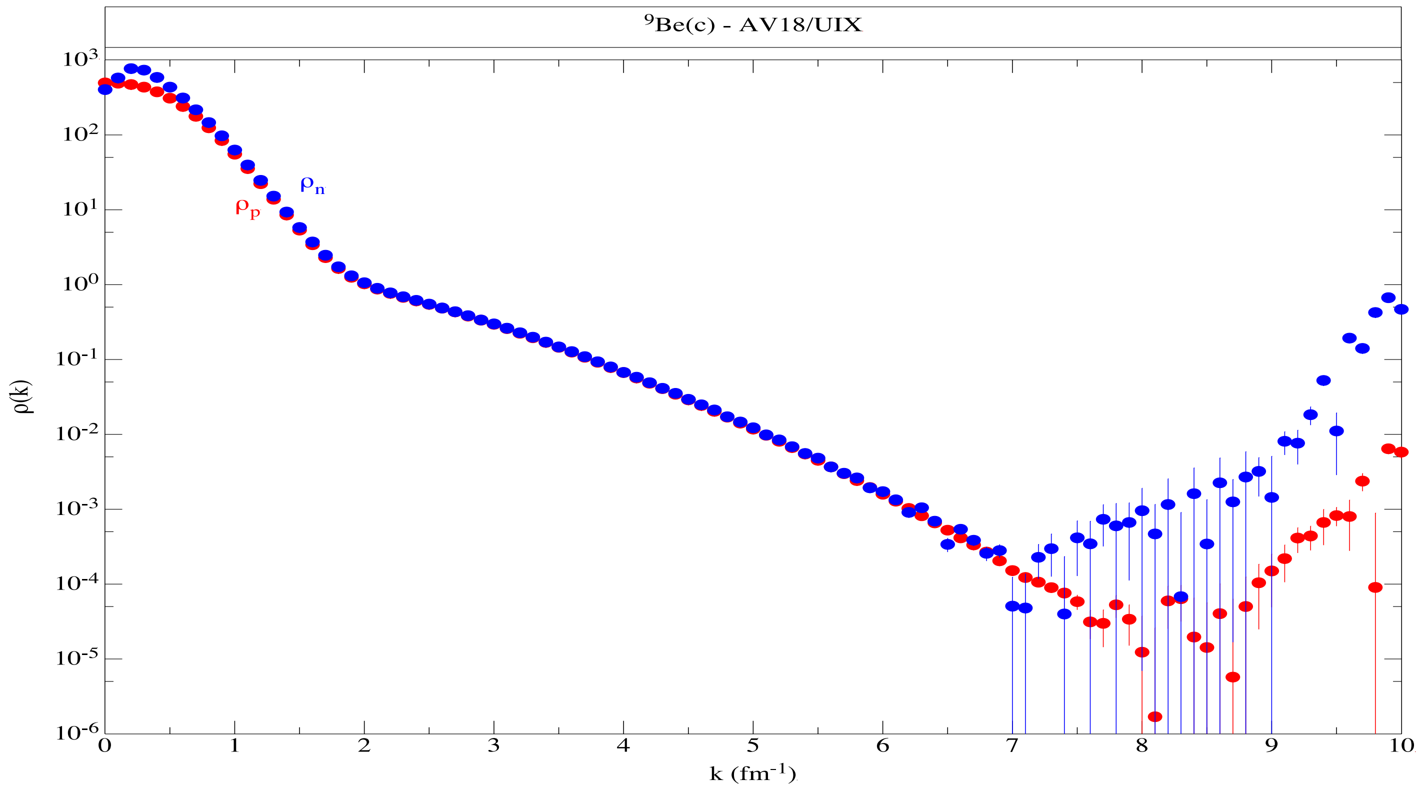
Realistic 3He Wave Function: Faddeev Equation



MS,PRC 2014

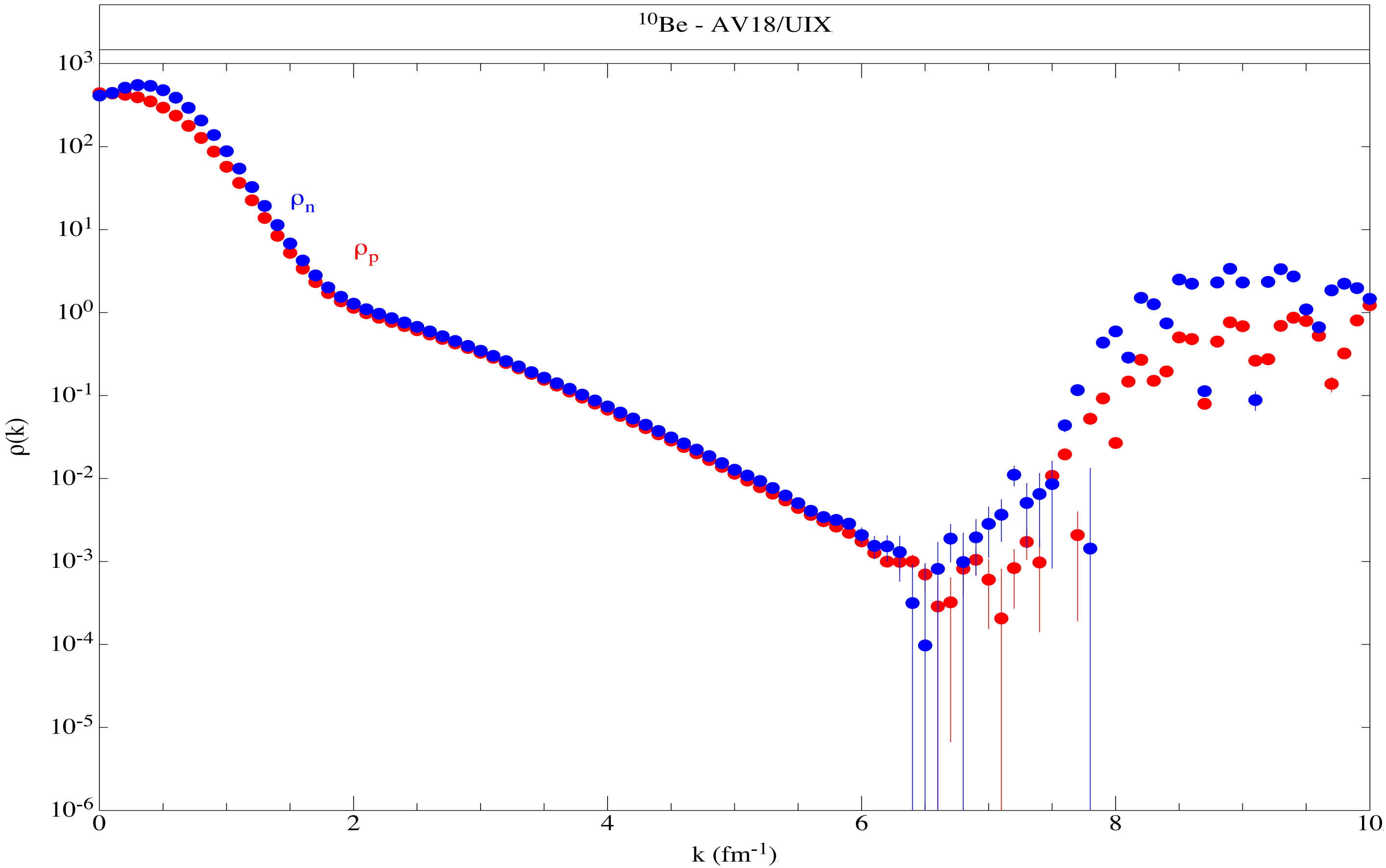
Be9 Variational Monte Carlo Calculation:

Robert Wiringa 2013 http://www.phy.anl.gov/theory/research/momenta/



B10 Variational Monte Carlo Calculation:

Robert Wiringa



Second Property: Inverse Fractional Dependence of High Momentum Component

$$x_p \cdot n_p^A(p) \approx x_n \cdot n_n^A(p) = \frac{a_{pn}}{2}(A,Z) n_d(p)$$

where $x_p = \frac{Z}{A}$ and $x_n = \frac{A-Z}{A}$.

$$n_{p/n}^A(p) \approx \frac{1}{2x_{p/n}} a_2(A,y) \cdot n_d(p)$$

Momentum distributions of p & n are inverse proportional to their fractions

Predictions: High Momentum Fractions

MS,arXiv:1210.3280,2012
Phys. Rev. C 2014

$$P_{p/n}(A,y) = \frac{1}{2x_{p/n}} a_2(A,y) \int_{k_F}^{\infty} n_d(p) d^3p$$

$$y = |x_p - x_n|$$

A	Pp(%)	Pn(%)
12	20	20
27	23	22
56	27	23
197	31	20

Predictions:

Minority Component has larger high momentum fraction:

MS,arXiv:1210.3280,2017
Phys. Rev. C 2014

Checking for He3

Energetic Neutron

$$E^p_{kin} = 14 \text{ MeV (p= 157 MeV/c)}$$

$$E^n_{kin} = 19 \text{ MeV (p= 182 MeV/c)}$$

Energetic Neutron
(Neff & Horiuchi)

$$E^p_{kin} = 13.97 \text{ MeV}$$

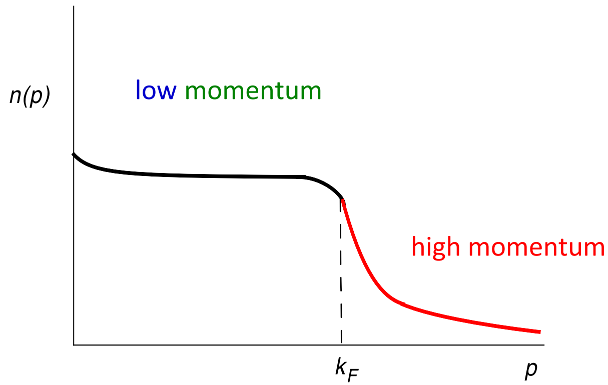
$$E^n_{kin} = 18.74 \text{ MeV}$$

VMC Estimates: Robert Wiringa

MS,arXiv:1210.3280
Phys. Rev. C 2014

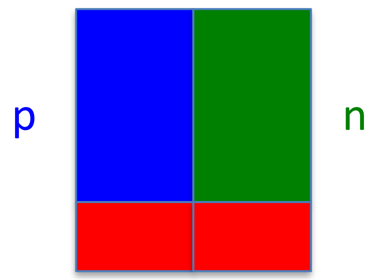
Table 1: Kinetic energies (in MeV) of proton and neutron

A	y	E^p_{kin}	E^n_{kin}	$E^p_{kin} - E^n_{kin}$
^{8}He	0.50	30.13	18.60	11.53
^{6}He	0.33	27.66	19.06	8.60
^{9}Li	0.33	31.39	24.91	6.48
^{3}He	0.33	14.71	19.35	-4.64
^{3}H	0.33	19.61	14.96	4.65
^{8}Li	0.25	28.95	23.98	4.97
^{10}Be	0.2	30.20	25.95	4.25
^{7}Li	0.14	26.88	24.54	2.34
^{9}Be	0.11	29.82	27.09	2.73
^{11}B	0.09	33.40	31.75	1.65



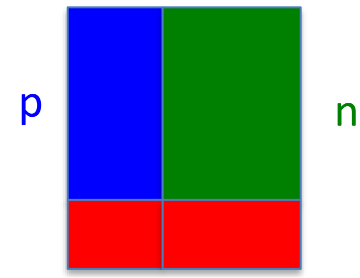
$$k_F = (3\pi^2 \rho_N)^{\frac{1}{3}}$$

Symmetric Nuclei



$$K_p = K_n$$

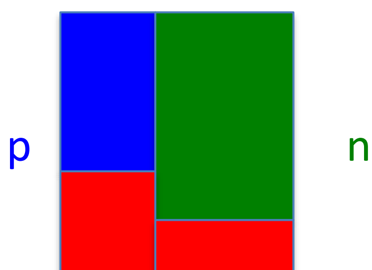
Asymmetric Nuclei



Conventional Theory: $K_n > K_p$

(Shell Model, HF, HO Ab Initio)

Asymmetric Nuclei



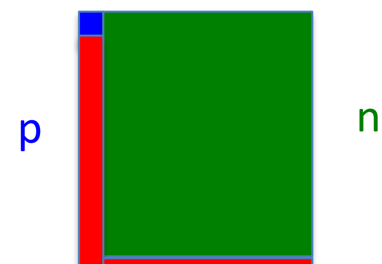
New Predictions

1. Per nucleon, more protons are in high momentum tail

2. Kin Energy Inversion

$K_p > K_n$?

Neutron Stars

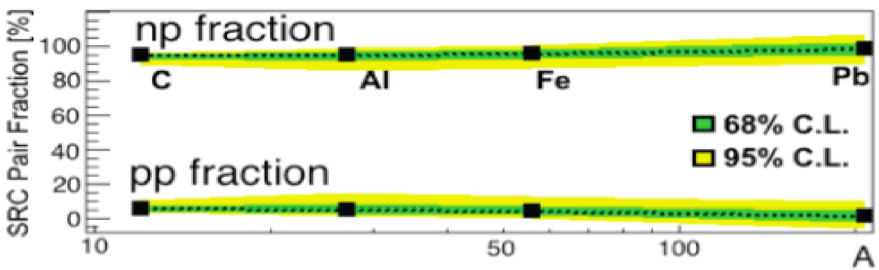


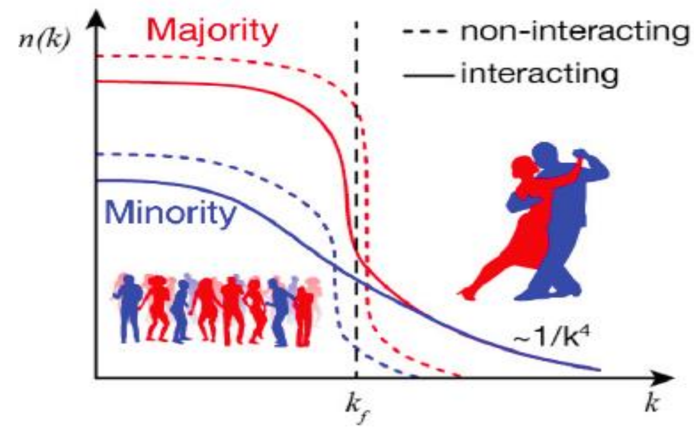
Protons my completely populate the high momentum tail

- Experimental Verification of Momentum Sharing Effects

- pn dominance persist for heavy nuclei

O. Hen, MS, L, Weinstein et.al. Science, 2014

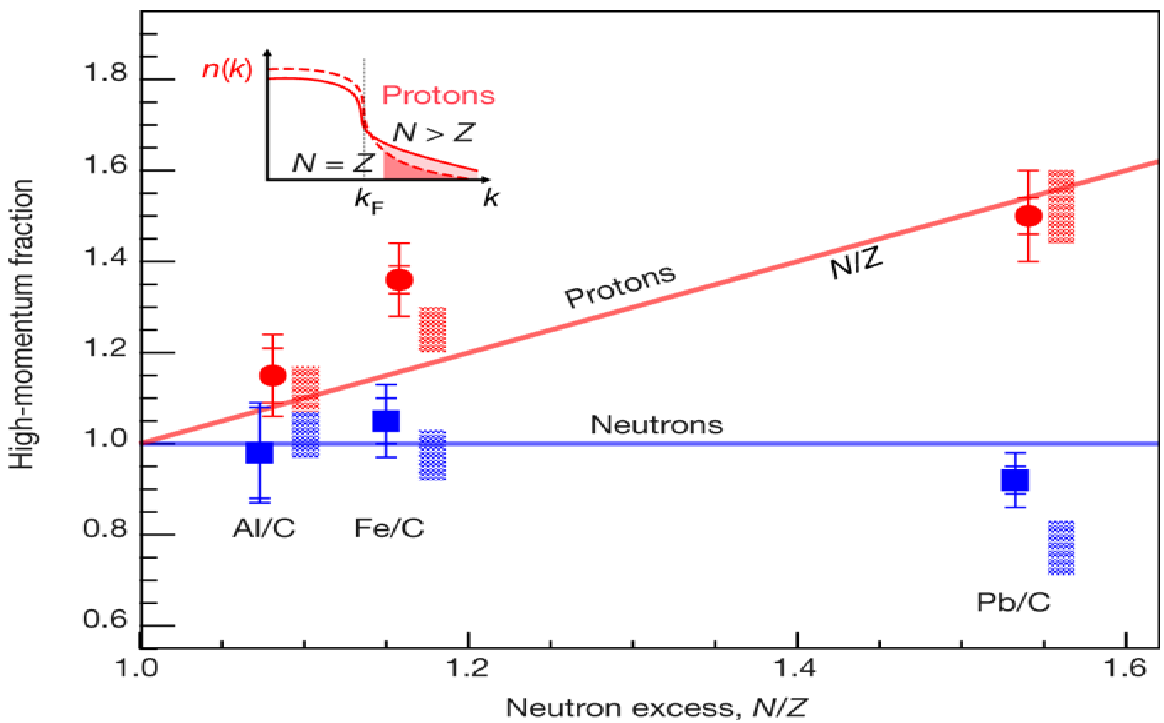




- verifying excess/suppression of high momentum protons/neutrons in neutron rich nuclei

- measuring proton and neutron momentum dependences separately

Duer et al, Nature 2018



- **New Properties of High Momentum Distribution of Nucleons in Asymmetric Nuclei**

  MS,arXiv:1210.3280.2012
  Phys. Rev. C 2014

- **Protons are more Energetic in Neutron Rich High Density Nuclear Matter**

  M. McGauley, MS
  arXiv:1102.3973,2011

- **First Experimental Indication**

  O. Hen, et.al.
  Science, 2014,

- **Confirmed by VMC calculations for A<12**

  R.B. Wiringa et al,
  Phys. Rev. C 2014

- **For Nuclear Matter**

  W. Dickhoff et al
  Phys. Rev. C 2014

- **For Medium/Heavy Nuclei**

  J. Ryckebusch, W.Cosyn
  M. Vanhalst., J.Phys 2015

- **In Light-Cone Approximation:**

  O.Artiles, M.S.
  Phys. Rev. C 2016

Implications/Predictions for Nuclear Physics and Astrophysics

- more/less protons/neutrons per nucleon in neutron rich nuclei
- protons are extremely energetic in Neutron Stars
- protons are more modified in neutron rich nuclei
- u-quarks are more modified than d-quarks in large A Nuclei

Experimental Implications:

- Flavor Dependence of EMC effect
- A dependence of NuTev Anomaly
- u/d modification can be checked in neutrino-nuclei or pvDIS processes

Implication for High Density Asymmetric Nuclear Matter

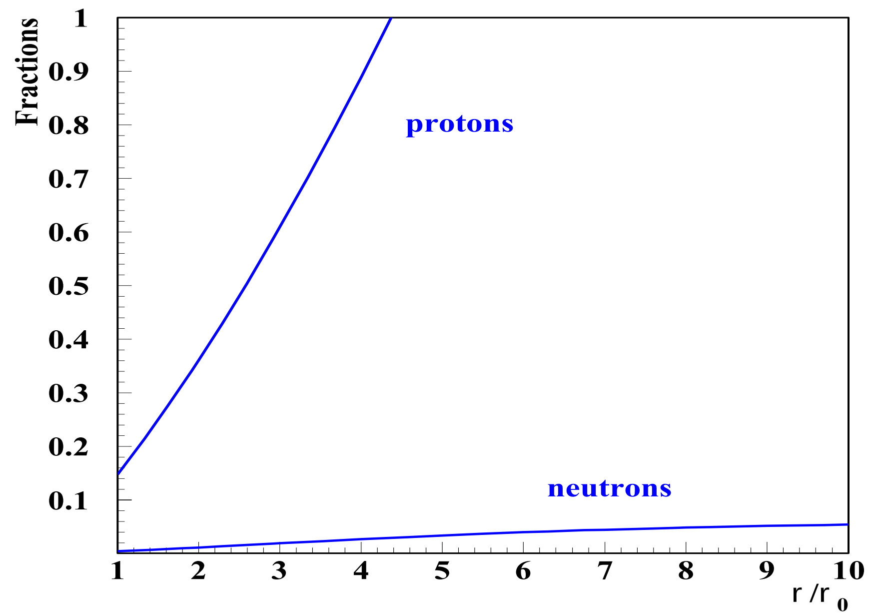
Implications: For Nuclear Matter

$$P_{p/n}(A,y) = \frac{1}{2x_{p/n}} a_2(A,y) \int_{k_F}^{\infty} n_d(p) d^3p$$

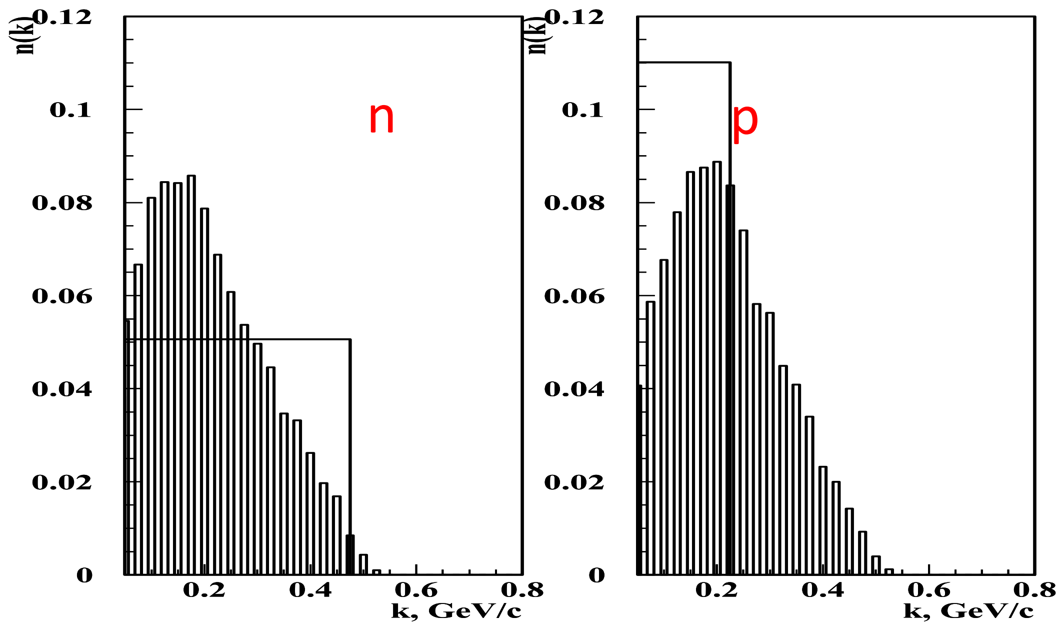
For $x_p = \frac{1}{9}$ and $y = \frac{7}{9}$
and using $k_{F,N} = (3\pi^2 x_N \rho)^{\frac{1}{3}}$

For $x_p = \frac{1}{9}$ and $y = \frac{7}{9}$
and using $k_{F,N} = (3\pi^2 x_N \rho)^{\frac{1}{3}}$

$$P_{p/n}(\rho, y) = \frac{a_2(\rho,y)}{2x_{p/n}} \int_{k_F} n_d(k) d^3k$$

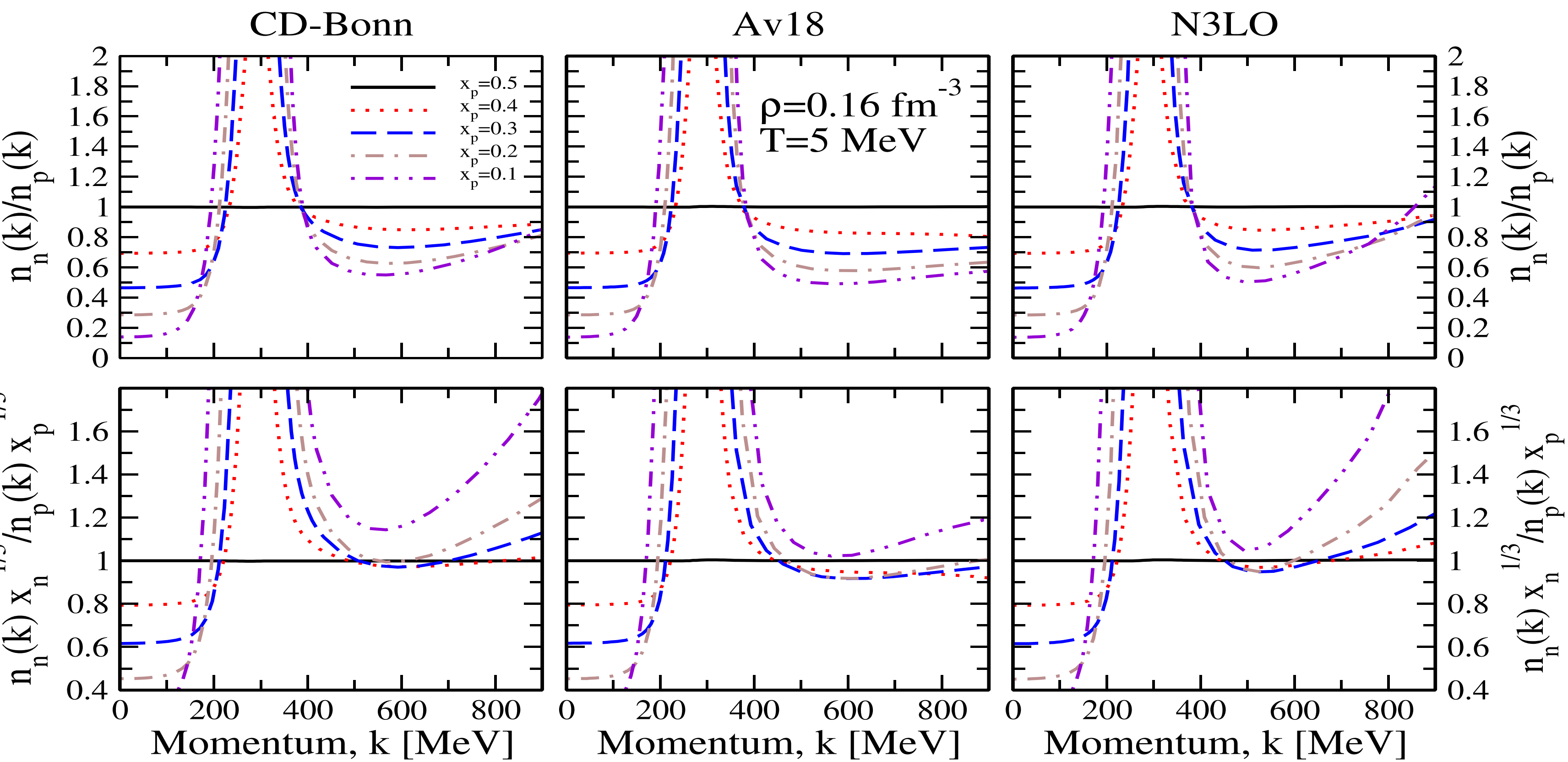


Implication for asymmetric nuclear matter



Asymmetric Nuclear Matter Calculations

A.Rios, A. Polls and W. H. Dickhoff, Phys. Rev. C 2014



Some Possible Implications of our Results

Cooling of Neutron Star:

Large concentration of the protons above the Fermi momentum will allow the condition for Direct URCA processes $p_p + p_e > p_n$ to be satisfied even if $x_p < \frac{1}{9}$. This will allow a situation in which intensive cooling of the neutron stars will be continued well beyond the critical point $x_p = \frac{1}{9}$.

Superfluidity of Protons in the Neutron Stars:

Transition of protons to the high momentum spectrum will smear out the energy gap which will remove the superfluidity condition for the protons. This will also result in significant changes in the mechanism of generation of neutron star magnetic fields.

Protons in the Neutron Star Cores:

The concentration of protons in the high momentum tail will result in proton densities $\rho_p \sim p_p^3 \gg k_{F,p}^3$. This will result in an equilibrium condition with "neutron skin" effect in which large concentration of protons will populate the core rather than the crust of the neutron star. This situation may provide very different dynamical conditions for generation of magnetic fields of the stars.

Isospin Locking and Large Masses of Neutron Stars

With an increase of the densities more and more protons move to the high momentum tail where they are in short range tensor correlations with neutrons. In this case on will expect that high density nuclear matter will be dominated by configurations with quantum numbers of tensor correlations $S = 1$ and $I = 0$. In such scenario protons and neutrons at large densities will be locked in the NN iso-singlet state. Such situation will double the threshold of inelastic excitation from $NN \to N\Delta$ to $NN \to \Delta\Delta(NN^*)$ transition thereby stiffening the equation of the state. This situation my explain the observed neutron star masses in Ref.[**?**] which are in agreement with the calculation of equation state that include only nucleonic degrees of freedom

Is the Observed Effects Universal for Two Component Asymmetric Fermi Systems?

- Start with Two Component Asymmetric Degenerate Fermi Gas

- Asymmetric: $\rho_1 << \rho_2$

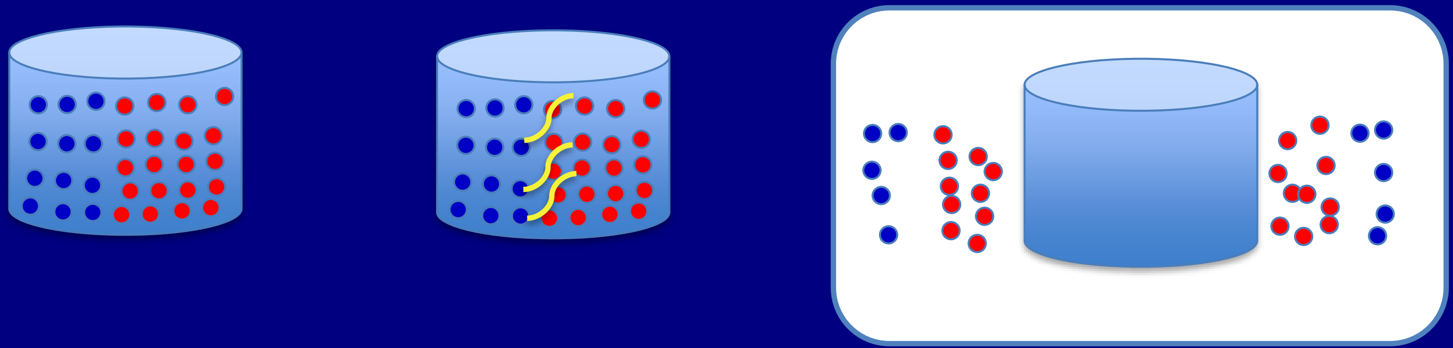
- Switch on the short-range interaction between two-component

- While interaction between each components is weak

- Spectrum of the small component gas will strongly deform

Cold Atoms

Possible Experiment with Trapped Atoms



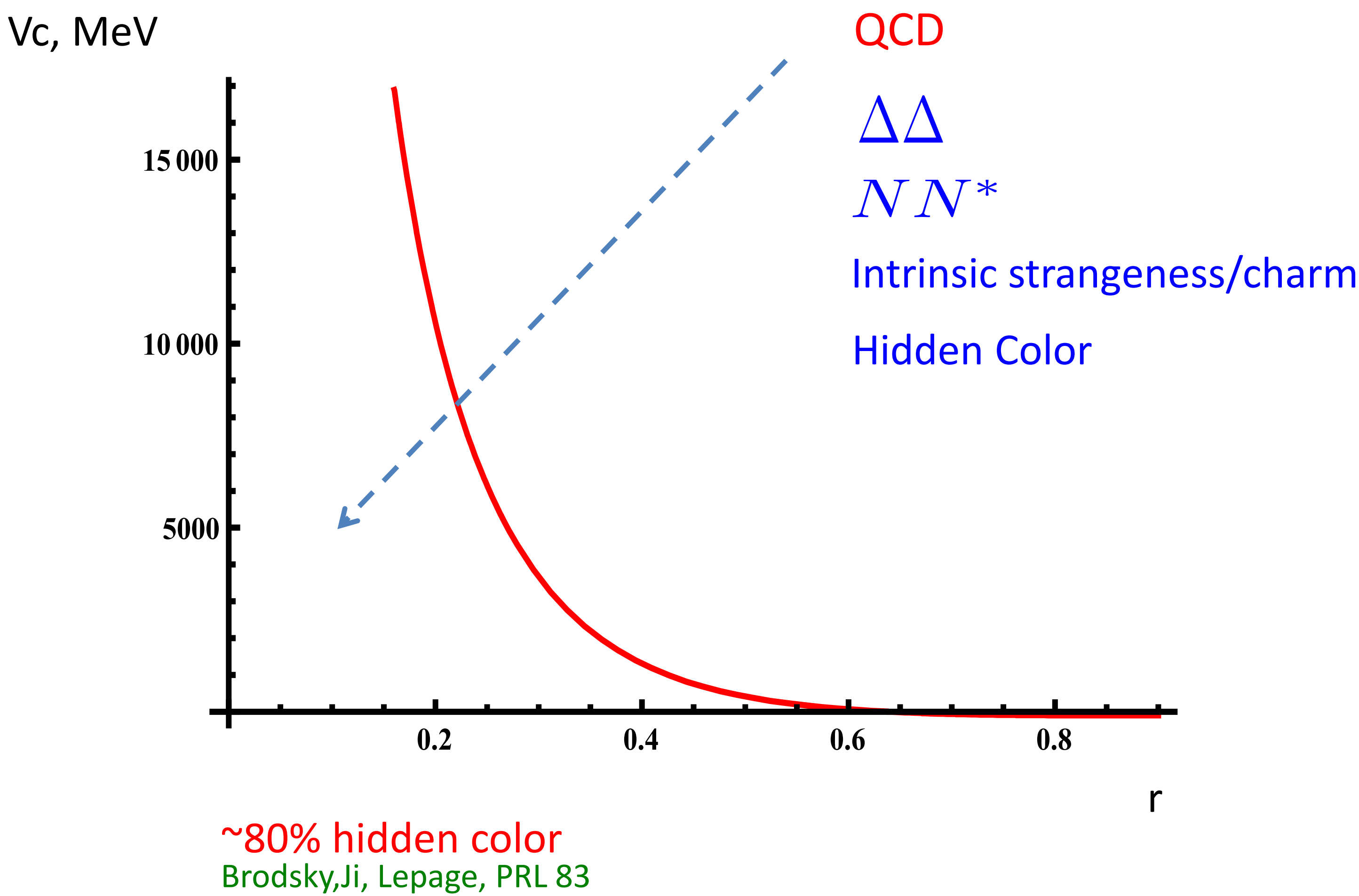
Conclusions and Outlook

- We observe two new properties of high momentum distribution of proton and neutron in nuclei

- Predicting more energetic/virtual protons in neutron reach matter

- May have strong implication for protons in neutron stars *Cooling & Magnetic Fields*

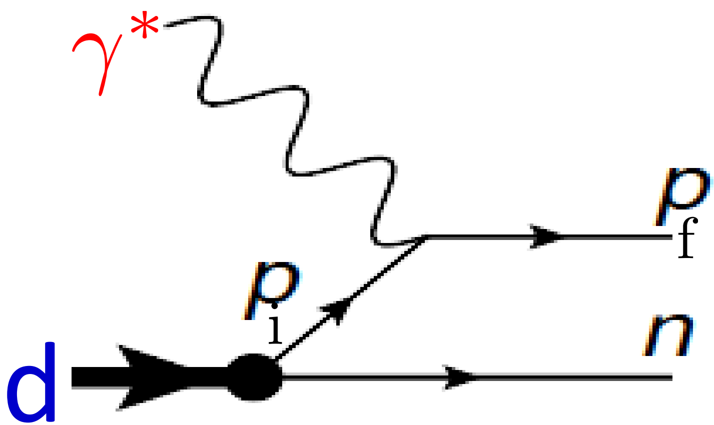
- Can be a universal effect for any two-component asymmetric fermi system with short-range interaction between unlike components


Vc, MeV
QCD
$\Delta\Delta$
NN^*
Intrinsic strangeness/charm
Hidden Color
15 000
10 000
5000
0.2
0.4
0.6
0.8
r
~80% hidden color
Brodsky,Ji, Lepage, PRL 83

Probing NN interaction at very short distances

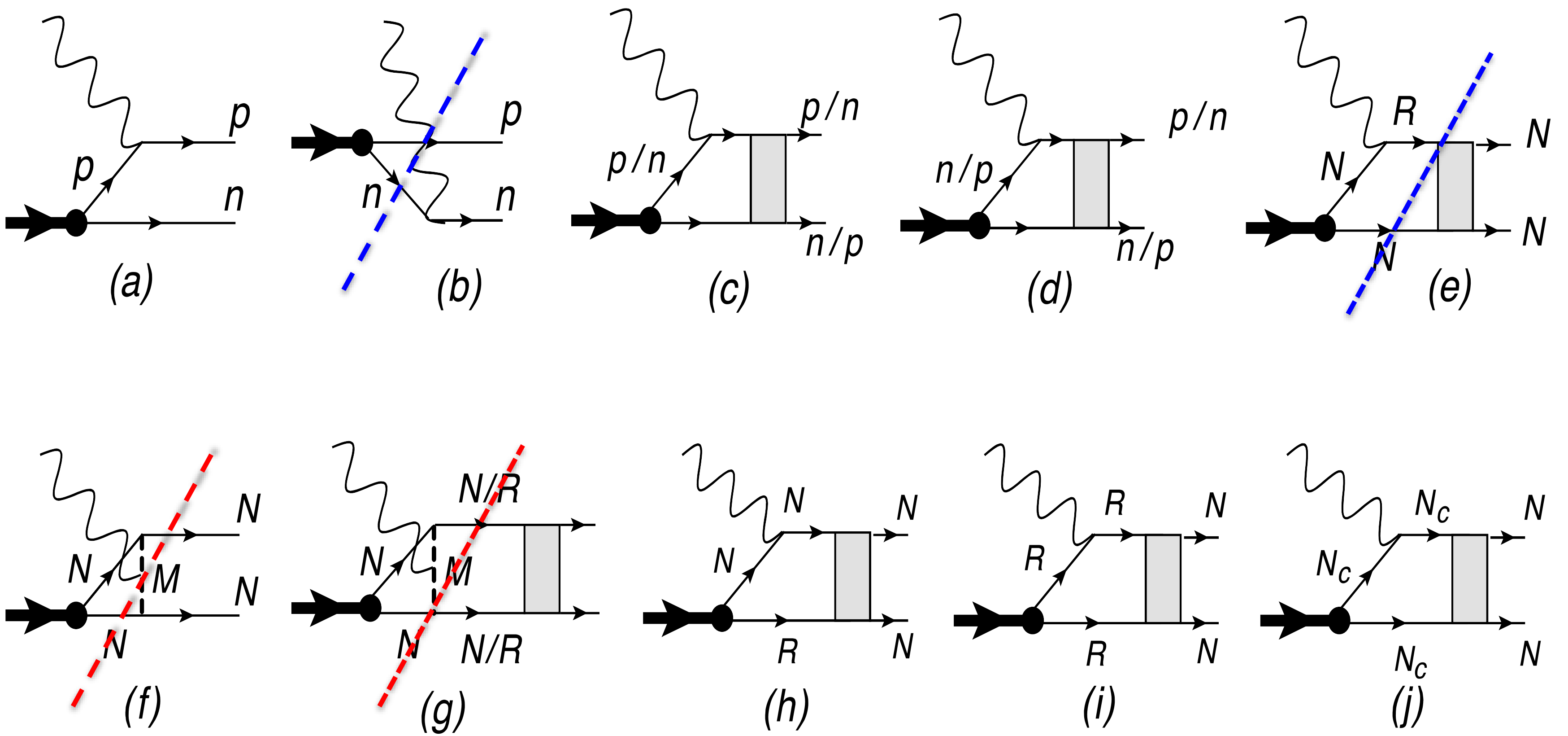
Considering reaction: $e + d \to e' + p_f + n$

$|p_i| = |p_f - q| \leq 550\ \mathrm{MeV/C}$



$|p_i| = |p_f - q| > 550\ \mathrm{MeV/C}$

Considering reaction: $e + d \to e' + p_f + n$



For the Deuteron it means, at Short Distances

$$\Psi_d = \Psi_{pn} + \Psi_{\Delta\Delta} + \Psi_{NN^*} + \Psi_{hc} \cdots + \Psi_{p\Lambda^0 K^0} \cdots$$

$$\Psi_{hc} = \Psi_{N_c, N_c}$$

The NN repulsive core can be due to the orthogonality of

$$\langle \Psi_{NonNucleonic} \mid \Psi_{NN} \rangle = 0$$

Some Paradigm Shift

Our current mindset about deuteron is fully non-relativistic, the observation that it has total spin, J=1 and parity, P=+, together with the relation that for non-relativistic wave function, P= $(-1)^l$, one concludes that the deuteron consists of S- and D- partial waves for proton-neutron system.

Paradigm Shift: The above reaction at high Q^2, measures the probability of observing proton and neutron in the deuteron at very large relative momenta.
In such a formulation the deuteron is not a composite system consisting of proton and neutron but it is a composite pseudo - vector (J=1, P=+) ``particle" from which one extracts proton and neutron.

How such a proton and neutron produced at such extremal conditions is related to the dynamical structure of Light-Front deuteron wave function, which may include internal elastic $pn \to pn$ as well as inelastic $\Delta\Delta \to pn$, $N^*N \to pn$ or $N_C N_C \to pn$ transitions.

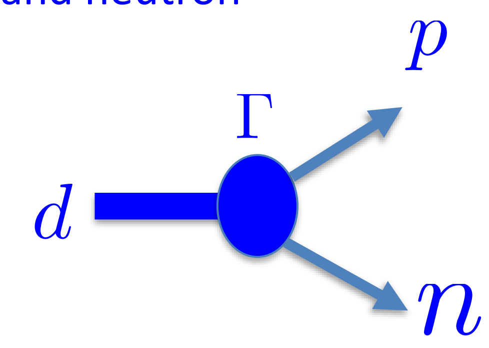
New Structure in the Deuteron and possible non-nucleonic components

M.S & Frank Vera PRL 2023

Paradigm shift:

- consider a deuteron not a nucleus that consist of proton and neutron
- but ***pseudovector composite particle*** from which we ***extract*** proton and neutron

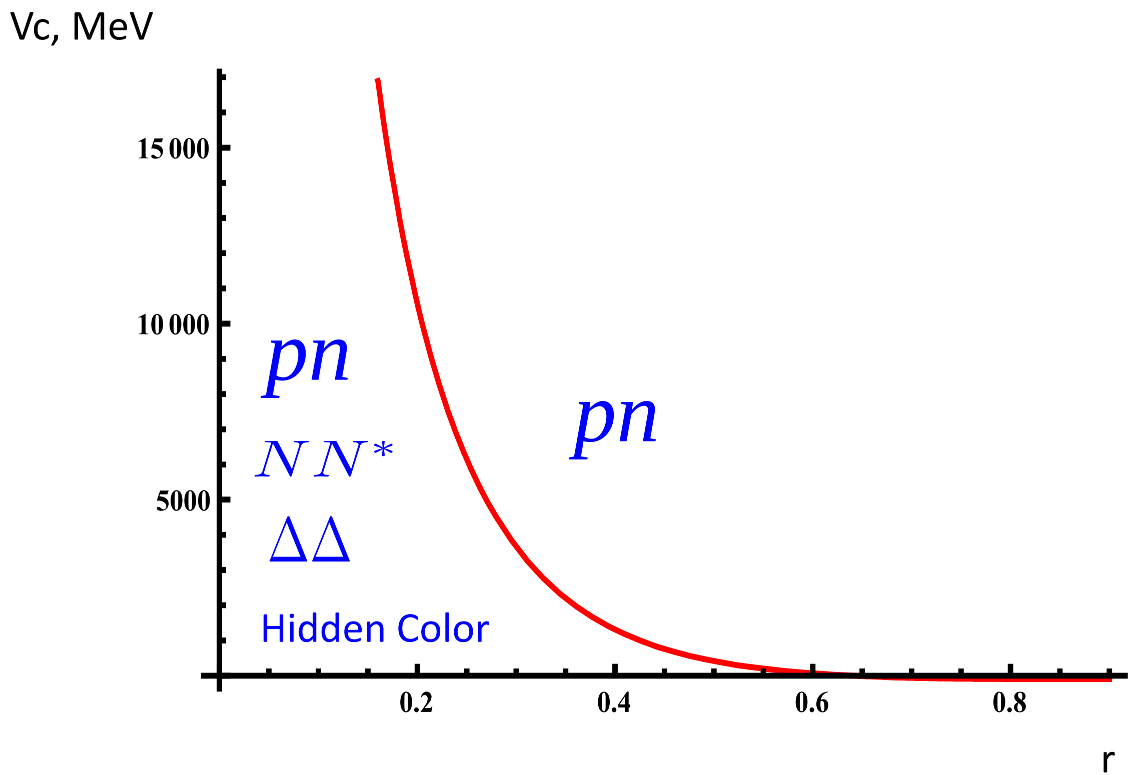
- Light-Front Deuteron wave function



$$\psi_d^{\lambda_d}(\alpha_i, p_\perp, \lambda_1\lambda_2) = -\frac{\bar{u}(p_2,\lambda_2)\bar{u}(p_1,\lambda_1)\Gamma_d\chi^{\lambda_d}}{\frac{1}{2}(m_d^2 - 4\frac{m_N^2+p_\perp^2}{\alpha_i(2-\alpha_i)})\sqrt{2(2\pi)^3}}$$

- Absorbing the energy denominator into the vertex function and using crossing symmetry

$$\psi_d^{\mu}(\alpha_i, p_\perp, \lambda_1, \lambda_2) = -\bar{u}(p_2,\lambda_2)\Gamma_d^\mu(k)\frac{(i\gamma_2\gamma_0)}{\sqrt{2}}\bar{u}(p_1,\lambda_1)^T = -\sum_{\lambda_1'}\bar{u}(p_1,\lambda_1)\Gamma_d^\mu\gamma_5\frac{\epsilon_{\lambda_1,\lambda_1'}}{\sqrt{2}}u(p_1,\lambda_1')$$



Vc, MeV
15 000
10 000
5000
pn
N N*
ΔΔ
Hidden Color
pn
0.2
0.4
0.6
0.8
r

$$\psi_d^\mu(\alpha_i, p_\perp, \lambda_1, \lambda_2) = -\bar{u}(p_2,\lambda_2)\Gamma_d^\mu(k)\frac{(i\gamma_2\gamma_0)}{\sqrt{2}}\bar{u}(p_1,\lambda_1)^T = -\sum_{\lambda_1'}\bar{u}(p_1,\lambda_1)\Gamma_d^\mu\gamma_5\frac{\epsilon_{\lambda_1,\lambda_1'}}{\sqrt{2}}u(p_1,\lambda_1')$$

- Γ_d^μ - is a four-vector, which can be constructed in a most general form satisfying time reversal, parity and charge conjugate symmetries

- Because the deuteron is a bound system, in addition to on-shell p_1 and p_2 four momenta one introduces

$$\Delta^\mu \equiv p_1^\mu + p_2^\mu - p_d^\mu \equiv (\Delta^-, \Delta^+, \Delta_\perp) = (\Delta^-, 0, 0)$$

$$\Delta^- = p_1^- + p_2^- - p_d^- = \frac{m_N^2 + k_\perp^2}{p_1^+} + \frac{m_N^2 + k_\perp^2}{p_2^+} - \frac{M_d^2}{p_d^+} = \frac{1}{p_d^+}\left[\frac{4(m_N^2 + k_\perp^2)}{\alpha_1(2-\alpha_1)} - M_d^2\right] = \frac{4}{p_d^+}\left[m_N^2 - \frac{M_d^2}{4} + k^2\right]$$

- Constructed vertex:

$$\Gamma_d^\mu = \Gamma_1\gamma^\mu + \Gamma_2\frac{(p_1 - p_2)^\mu}{2m_N} + \Gamma_3\frac{\Delta^\mu}{2m_N} + \Gamma_4\frac{(p_1 - p_2)^\mu \not{\Delta}}{4m_N^2}$$
$$+ i\Gamma_5\frac{1}{4m_N^3}\gamma_5\epsilon^{\mu\nu\rho\gamma}(p_d)_\nu(p_1 - p_2)_\rho(\Delta)_\gamma + \Gamma_6\frac{\Delta^\mu \not{\Delta}}{4m_N^2}$$

High Momentum Transfer Kinematics

For large Q^2 limit, Light-Front momenta for the reaction are chosen as follows:

$$p_d^\mu \equiv (p_d^-, p_d^+, p_{d\perp}) = \left(\frac{Q^2}{x\sqrt{s}} \left[1 + \frac{x}{\tau} - \sqrt{1 + \frac{x^2}{\tau}} \right], \frac{Q^2}{x\sqrt{s}} \left[1 + \frac{x}{\tau} + \sqrt{1 + \frac{x^2}{\tau}} \right], 0_\perp \right)$$

$$q^\mu \equiv (q^-, q^+, q_\perp) = \left(\frac{Q^2}{x\sqrt{s}} \left[1 - x + \sqrt{1 + \frac{x^2}{\tau}} \right], \frac{Q^2}{x\sqrt{s}} \left[1 - x - \sqrt{1 + \frac{x^2}{\tau}} \right], 0_\perp \right)$$

where $s = (q + p_d)^2$, $\tau = \frac{Q^2}{M_d^2}$ and $x = \frac{Q^2}{M_d q_0}$, with q_0 being virtual photon energy in the deuteron rest frame.

- One observes that for fixed x, $p_d^+ \sim \sqrt{Q2} \gg m_N$

$$\Delta^\mu \equiv p_1^\mu + p_2^\mu - p_d^\mu \equiv (\Delta^-, \Delta^+, \Delta_\perp) = (\Delta^-, 0, 0),$$

where

$$\Delta^- = p_1^- + p_2^- - p_d^- = \frac{m_N^2 + k_\perp^2}{p_1^+} + \frac{m_N^2 + k_\perp^2}{p_2^+} - \frac{M_d^2}{p_d^+}$$

$$= \frac{1}{p_d^+}\left[\frac{4(m_N^2 + k_\perp^2)}{\alpha_1(2-\alpha_1)} - M_d^2\right] = \frac{4}{p_d^+}\left[m_N^2 - \frac{M_d^2}{4} + k^2\right].$$

In high Q^2 limit $\quad \frac{\Delta^-}{2m_N} \ll 1$

$$\Gamma_d^\mu = \Gamma_1 \gamma^\mu + \Gamma_2 \frac{(p_1 - p_2)^\mu}{2m_N} + \cancel{\Gamma_3 \frac{\Delta^\mu}{2m_N}} + \cancel{\Gamma_4 \frac{(p_1 - p_2)^\mu \not\Delta}{4m_N^2}}$$

$$+ i\Gamma_5 \frac{1}{4m_N^3} \gamma_5 \epsilon^{\mu\nu\rho\gamma} (p_d)_\nu (p_1 - p_2)_\rho (\Delta)_\gamma + \cancel{\Gamma_6 \frac{\Delta^\mu \not\Delta}{4m_N^2}}$$

Consider: $\epsilon^{\mu,+,\perp,-} p_{d,-} k_\perp \Delta_+$

Since: $p_{d,-} = \frac{1}{2} p_d^+$ and $\Delta_+ = \frac{1}{2}\Delta^-$ then $p_d^+ \Delta^- = p_d^+ \frac{1}{p_d^+}\left[\frac{4(m_N^2+k_\perp^2)}{\alpha_1(2-\alpha_1)} - M_d^2\right] = \left[\frac{4(m_N^2+k_\perp^2)}{\alpha_1(2-\alpha_1)} - M_d^2\right]$

$\epsilon^{\mu,+,\perp,-} p_{d,-} k_\perp \Delta_+ = \frac{1}{4}\epsilon^{\mu,+,\perp,-} p_d^+ k_\perp \Delta^-$ Leading Oder!

$$\Gamma_d^\mu = \Gamma_1 \gamma^\mu + \Gamma_2 \frac{(p_1 - p_2)^\mu}{2m_N} + \Gamma_3 \frac{\Delta^\mu}{2m_N} + \Gamma_4 \frac{(p_1 - p_2)^\mu \not\Delta}{4m_N^2}$$
$$+ i\Gamma_5 \frac{1}{4m_N^3}\gamma_5 \epsilon^{\mu\nu\rho\gamma}(p_d)_\nu (p_1 - p_2)_\rho (\Delta)_\gamma + \Gamma_6 \frac{\Delta^\mu \not\Delta}{4m_N^2}$$

$$\psi_d^{\lambda_d}(\alpha_i, k_\perp) = -\sum_{\lambda_2,\lambda_1,\lambda_1'} \bar{u}(-k,\lambda_2)\left\{\Gamma_1\gamma^\mu + \Gamma_2 \frac{\tilde{k}^\mu}{m_N} + \sum_{i=1}^{2} i\Gamma_5 \frac{1}{8m_N^3}\epsilon^{\mu+i-}p_d'^{+} k_i \Delta'^{-}\right\}\gamma_5 \frac{\epsilon_{\lambda_1,\lambda_i'}}{\sqrt{2}} u(k,\lambda_1') s_\mu^{\lambda_d}$$

where $\tilde{k}^\mu = (0, k_z, k_\perp)$

$$\psi_d^{\lambda_d}(\alpha_i, k_\perp) = -\sum_{\lambda_2,\lambda_1,\lambda_1'} \bar{u}(-k,\lambda_2)\left\{\Gamma_1\gamma^\mu + \Gamma_2\frac{\tilde{k}^\mu}{m_N} + \sum_{i=1}^{2} i\Gamma_5\frac{1}{8m_N^3}\epsilon^{\mu+i-}p_d'^{+}k_i\Delta'^{-}\right\}\gamma_5\frac{\epsilon_{\lambda_1,\lambda_i'}}{\sqrt{2}}u(k,\lambda_1')s_\mu^{\lambda_d}$$

$$\psi_d^{\lambda_d}(\alpha_1, k_t, \lambda_1, \lambda_2) = \sum_{\lambda_1'} \phi_{\lambda_2}^\dagger \sqrt{E_k}\left[\frac{U(k)}{\sqrt{4\pi}}\sigma\mathbf{s}_\mathbf{d}^{\lambda_\mathbf{d}} - \frac{W(k)}{\sqrt{4\pi}\sqrt{2}}\left(\frac{3(\sigma\mathbf{k})(\mathbf{k}\mathbf{s}_\mathbf{d}^\lambda)}{k^2} - \sigma\mathbf{s}_\mathbf{d}^\lambda\right) + (-1)^{\frac{1+\lambda_d}{2}}P(k)Y_1^{\lambda_d}(\theta,\phi)\delta^{1,|\lambda_d|}\right]\frac{\epsilon_{\lambda_1,\lambda_1'}}{\sqrt{2}}\phi_{\lambda_1'}$$

$$U(k) = \frac{2\sqrt{4\pi}\sqrt{E_k}}{3}\left[\Gamma_1(2+\frac{m_N}{E_k}) + \Gamma_2\frac{k^2}{m_N E_k}\right]$$

$$W(k) = \frac{2\sqrt{4\pi}\sqrt{2E_k}}{3}\left[\Gamma_1(1-\frac{m_N}{E_k}) - \Gamma_2\frac{k^2}{m_N E_k}\right]$$

Where: $Y_1^{\pm}(\theta,\phi) = \mp i\sqrt{\frac{3}{4\pi}}\sum_{i=1}^{2}\frac{(k\times s_d^{\pm 1})_z}{k}$

$$P(k) = \sqrt{4\pi}\frac{\Gamma_5(k)\sqrt{E_k}}{\sqrt{3}}\frac{k^3}{m_N^3}$$

fully relativistic: in addition to $\frac{k^{l=1}}{m_N}$ term

has additional $\frac{k^2}{m_N^2}$ term

Light Front Density Matrix and Momentum Distribution

$$\psi_d^{\lambda_d}(\alpha_1, k_t, \lambda_1, \lambda_2) = \sum_{\lambda_1'} \phi^\dagger_{\lambda_2} \sqrt{E_k} \left[\frac{U(k)}{\sqrt{4\pi}} \sigma \mathbf{s}_\mathbf{d}^{\lambda_\mathbf{d}} - \frac{W(k)}{\sqrt{4\pi}\sqrt{2}} \left(\frac{3(\sigma\mathbf{k})(\mathbf{k}\mathbf{s}_\mathbf{d}^\lambda)}{k^2} - \sigma \mathbf{s}_\mathbf{d}^\lambda \right) + (-1)^{\frac{1+\lambda_d}{2}} P(k) Y_1^{\lambda_d}(\theta,\phi) \delta^{1,|\lambda_d|} \right] \frac{\epsilon_{\lambda_1,\lambda_1'}}{\sqrt{2}} \phi_{\lambda_1'}$$

$$\rho_d(\alpha, k_\perp) = \frac{n_d(k,k_\perp)}{2-\alpha}$$

$$n_d(k,k_\perp) = \frac{1}{3} \sum_{\lambda_d=-1}^{1} \mid \psi_d^{\lambda_d}(\alpha, k_\perp) \mid^2 = \frac{1}{4\pi} \left(U(k)^2 + W(k)^2 + \frac{k_\perp^2}{k^2} P^2(k) \right)$$

Baryonic and Momentum Sum Rules $\int \rho_d(\alpha, k_\perp) \frac{d\alpha}{\alpha} = 1 \quad \text{and} \quad \int \alpha \rho_d(\alpha, k_\perp) \frac{d\alpha}{\alpha} = 1$

$$\int \left(U(k)^2 + W(k)^2 + \frac{2}{3} P^2(k) \right) k^2 dk = 1.$$

Non-Nucleonic Components and the New Structure

$$n_d(k,k_\perp) = \frac{1}{3}\sum_{\lambda_d=-1}^{1} \mid \psi_d^{\lambda_d}(\alpha,k_\perp)\mid^2 = \frac{1}{4\pi}\left(U(k)^2 + W(k)^2 + \frac{k_\perp^2}{k^2}P^2(k)\right)$$

- Momentum distribution depends on $k_\perp$ separately

- *This is impossible for non-relativistic quantum mechanics of the deuteron since in this case the potential of the interaction is real (no inelasticities) and the solution of Lippmann-Schwinger (or Schroedinger) equation for partial S- and D-waves satisfies ``angular condition", according to which the momentum distribution in unpolarized deuteron depends on the magnitude of relative momentum only.*

- On the other hand, in the relativistic domain the definition of the interaction potential is not straightforward to allow to use quantum-mechanical arguments in claiming that momentum distribution should satisfy the angular condition (i.e. depends on magnitude of k only).

- However, for the Light-Front, there is a remarkable theorem (Frankfurt, Mankiewicz, Sawitzky, Strikman, 1990) which states that if one considers only pn component in the deuteron, then for most acceptable forms of NN potential – constructed from elastic pn → pn scattering, the angular condition should be satisfied also for LF momentum distribution.

$$T_{NN}(\alpha_i, k_{i\perp}, \alpha_f, k_{f,\perp}) \equiv T_{NN}(k_{i,z}, k_{i\perp}, k_{f,z}, k_{f,\perp}) = V(k_{i,z}, k_{i\perp}, k_{f,z}, k_{f,\perp}) + \int V(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) \times \frac{d^3 k_m}{(2\pi)^3\sqrt{m^2 + k_m^2}} \frac{T_{NN}(k_{m,z}, k_{m\perp}, k_{f,z}, k_{f,\perp})}{4(k_m^2 - k_f^2)}$$

-The realization of the angular condition for relativistic case will require that light-front potential to satisfy a condition

$$V(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) = V(\vec{k}_i^2, (\vec{k}_m - \vec{k}_i)^2)$$

- Lorentz invariance for on-shell NN amplitude requires

$$T_{NN}^{on\ shell}(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) = T_{NN}^{on\ shell}(\vec{k}_i^2, (\vec{k}_m - \vec{k}_i)^2)$$

- Existence of the Born term indicates that

$$T_{NN}^{on\ shell}(\vec{k}_i^2, (\vec{k}_m - \vec{k}_i)^2) = V_{NN}^{on\ shell}(\vec{k}_i^2, (\vec{k}_m - \vec{k}_i)^2) + \int V_{NN}(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) \times \frac{d^3 k_m}{(2\pi)^3 \sqrt{m^2 + k_m^2}} \frac{T_{NN}(k_{m,z}, k_{m\perp}, k_{f,z}, k_{f,\perp})}{4(k_m^2 - k_f^2)}$$

- Iterating the equation around the on-shell kinematic point.

$$T_{NN}(\alpha_i, k_{i\perp}, \alpha_f, k_{f,\perp}) \equiv T_{NN}(k_{i,z}, k_{i\perp}, k_{f,z}, k_{f,\perp}) = V(k_{i,z}, k_{i\perp}, k_{f,z}, k_{f,\perp}) + \int V(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) \times \frac{d^3 k_m}{(2\pi)^3 \sqrt{m^2 + k_m^2}} \frac{T_{NN}(k_{m,z}, k_{m\perp}, k_{f,z}, k_{f,\perp})}{4(k_m^2 - k_f^2)}$$

- will result in:

$$T_{NN}(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) = T_{NN}(\vec{k}_i^2, (\vec{k}_m - \vec{k}_i)^2)$$

$$V_{NN}(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) = V_{NN}(\vec{k}_i^2, (\vec{k}_m - \vec{k}_i)^2)$$

for the general case

- V_{NN} – analytic function of angular momentum and it does not diverge exponentially in the complex-angular momentum space it was shown that also for the off-shell case

- For Non-nucleonic components no such iteration can be done

$$T_{NN}(k_{i,z}, k_{i\perp}, k_{f,z}, k_{f,\perp}) = \int V_{NN^*}(k_{i,z}, k_{i\perp}, k_{m,z}, k_{m,\perp}) \frac{d^3 k_m}{(2\pi)^3 \sqrt{m_m^2 + k_m^2}} \frac{T_{N^*N}(k_{m,z}, k_{m\perp}, k_{f,z}, k_{f,\perp})}{4(k_m^2 - k_f^2 + m_m^2 - m_N^2)}$$

- transition amplitudes such is $T_{\Delta\Delta\to NN}$, $T_{N^*,N\to NN}$ or $T_{N_c,N_c\to NN}$ where $N^c N^c$ represents a hidden color component in the deuteron could not be described with any combination of interaction potentials that satisfies angular condition

- if Γ_5 term is not zero then it should originate from non-nucleonic component in the deuteron.

- Our prediction is that the observation of LF momentum distribution depending on the center of mass k and $k_\perp$ separately will indicate the presence of non-nucleonic component in the deuteron

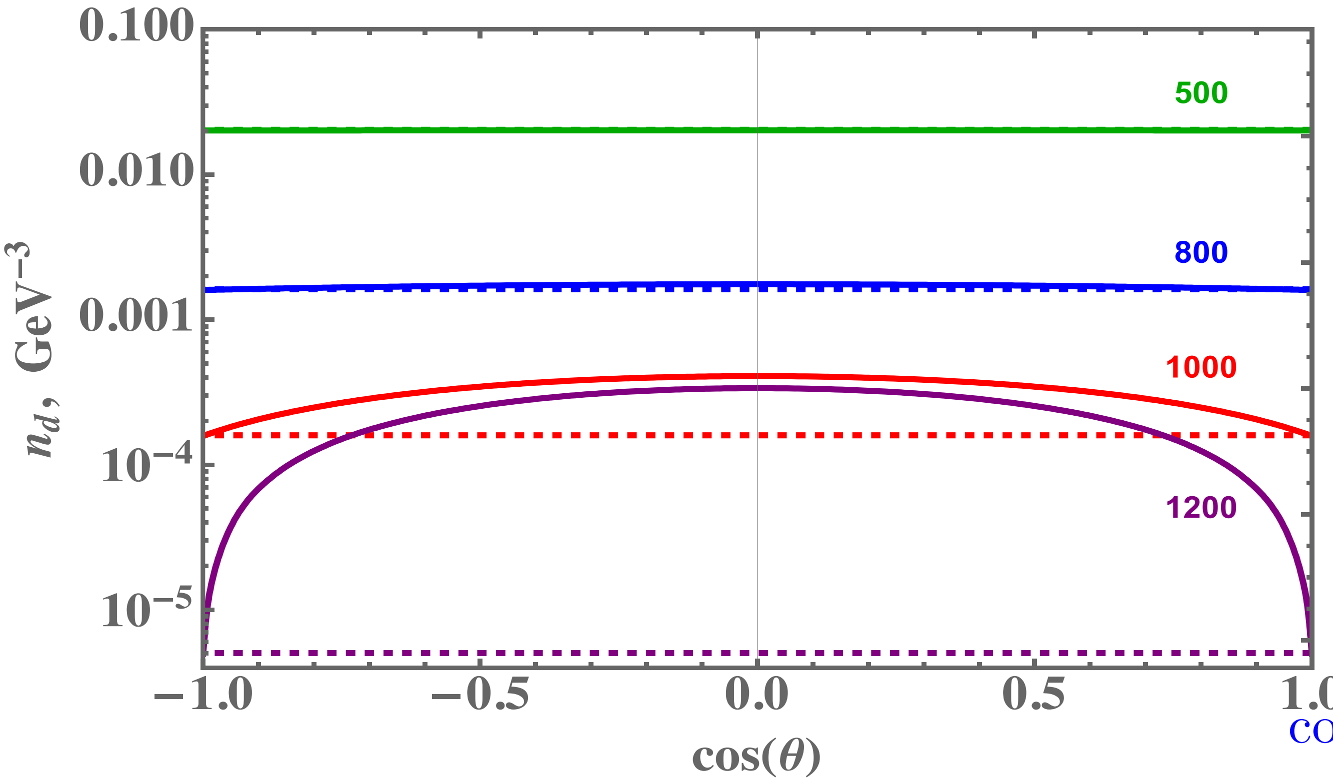
Estimate of the effect

$$n_d(k,k_\perp) = \frac{1}{3}\sum_{\lambda_d=-1}^{1} |\,\psi_d^{\lambda_d}(\alpha,k_\perp)\,|^2 = \frac{1}{4\pi}\left(U(k)^2 + W(k)^2 + \frac{k_\perp^2}{k^2}P^2(k)\right)$$

$$P(k) = \sqrt{4\pi}\,\frac{\Gamma_5(k)\sqrt{E_k}}{\sqrt{3}}\,\frac{k^3}{m_N^3}$$

$$\Gamma_5(k) = \frac{A}{(1+\frac{k^2}{0.71})^2}$$

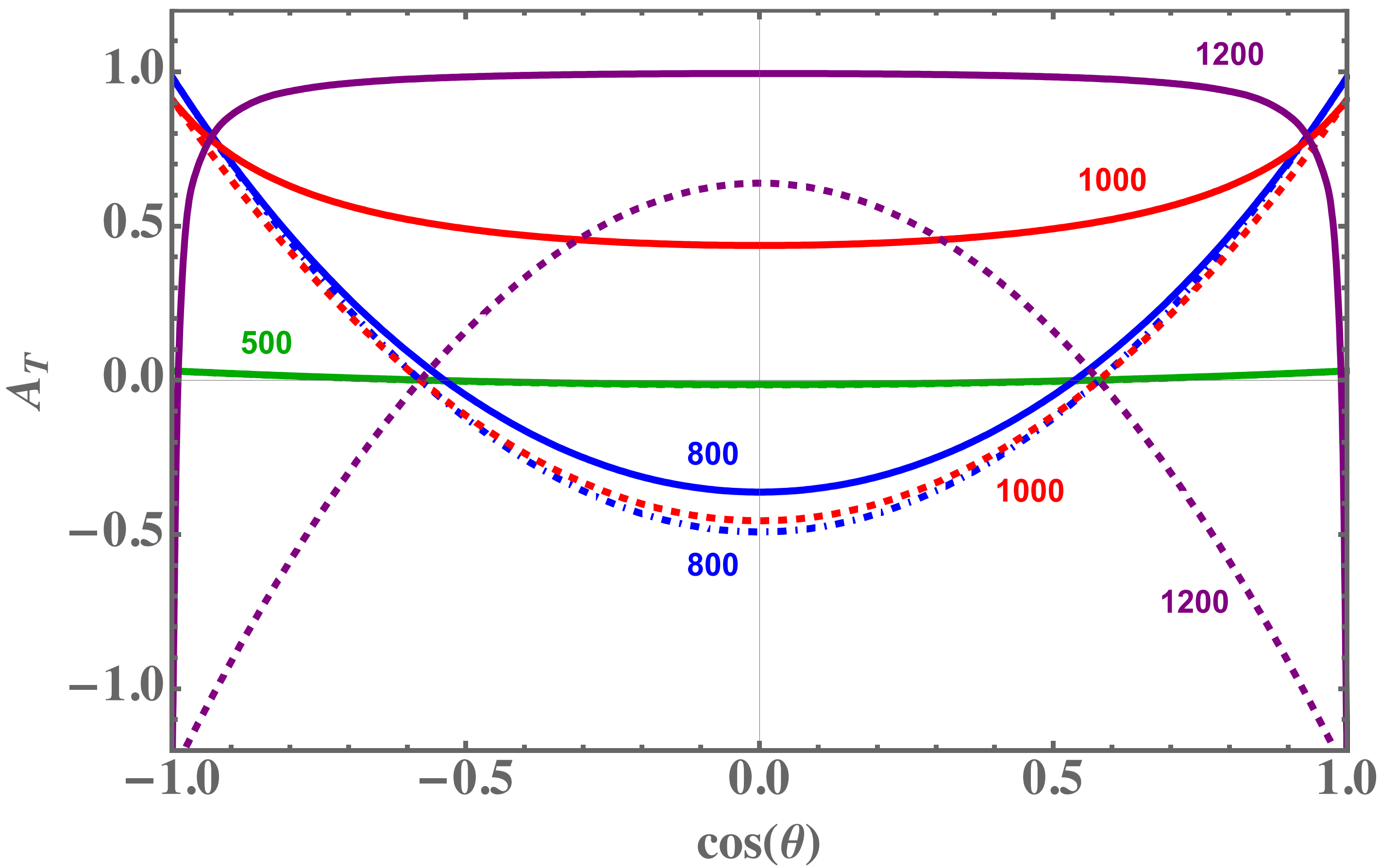
A is estimated by assuming 1% contribution to the total normalization.



$$\cos\theta = \frac{(\alpha-1)E_k}{k}$$

Estimate of the effect

$$A_T = \frac{n_d^{\lambda_d=1}(k,k_\perp)+n_d^{\lambda_d=-1}(k,k_\perp)-2n_d^{\lambda_d=0}(k,k_\perp)}{n_d(k,k_\perp)}$$

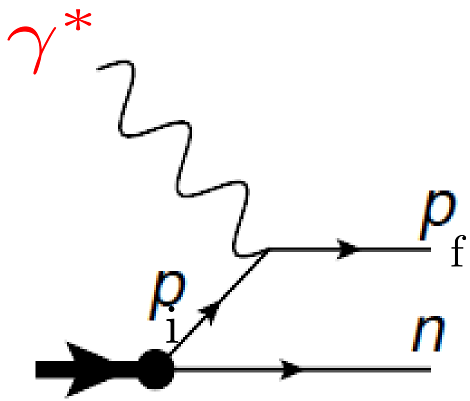


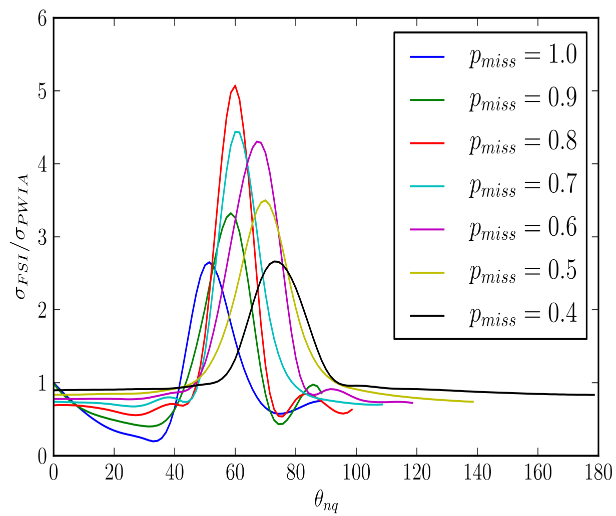
$$\cos\theta = \frac{(\alpha-1)E_k}{k}$$

Possibility of Experimental Verification

Considering reaction: $e + d \to e' + p_f + n$

$$|p_i| = |p_f - q| \gtrsim 800\text{MeV}/c$$





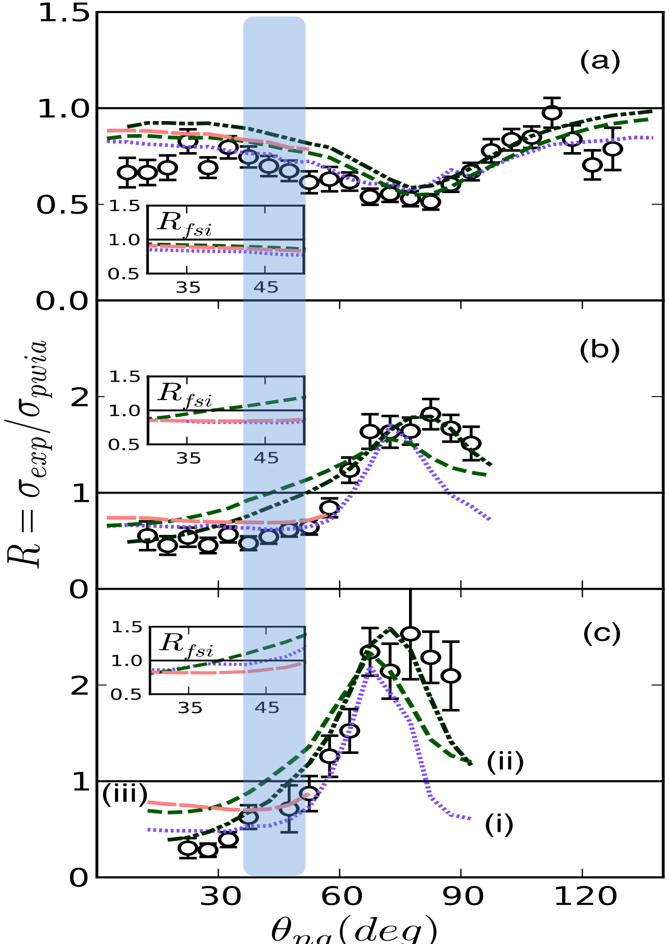
PAC-36, 2010

■ **E12-10-003 (p_m □ 300 MeV)**: "Deuteron Electro-Disintegration at Very High Missing Momentum"

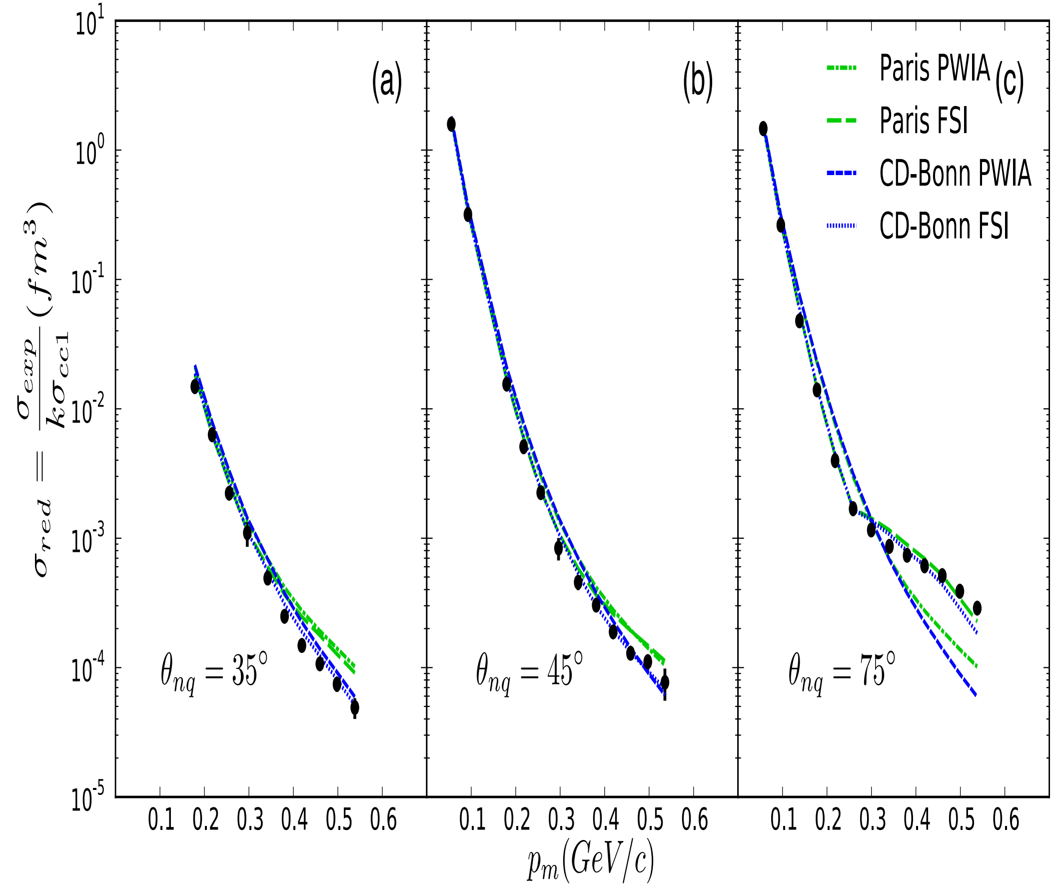
Rating: B+

data are essential to constrain further theory developments. Overall the experiment was viewed very highly; the lower rating simply reflects the likelihood that the data will not reveal any particular surprise and that their impact may thus be limited to experts in the field.

Probing Deuteron at Small Distances at large Q^2



JLab, $Q^2 = 3.5\ GeV^2$

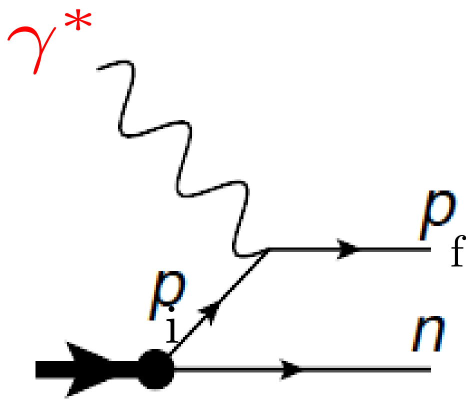


Boeglin et al PRL 2011, deuteron probed at up to 550MeV/c

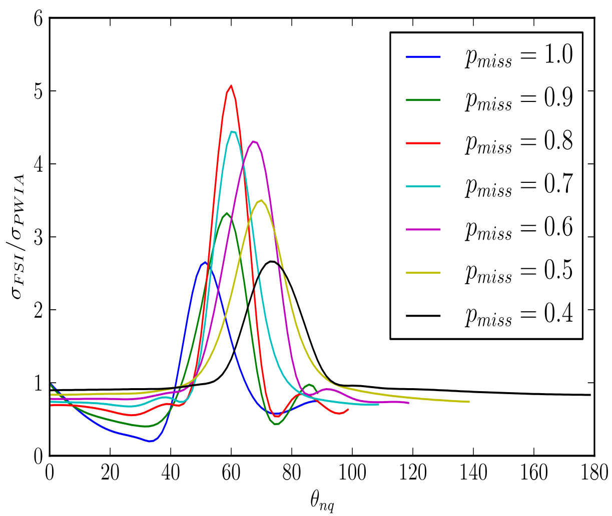
Possibility of Experimental Verification

Considering reaction: $e + d \to e' + p_f + n$

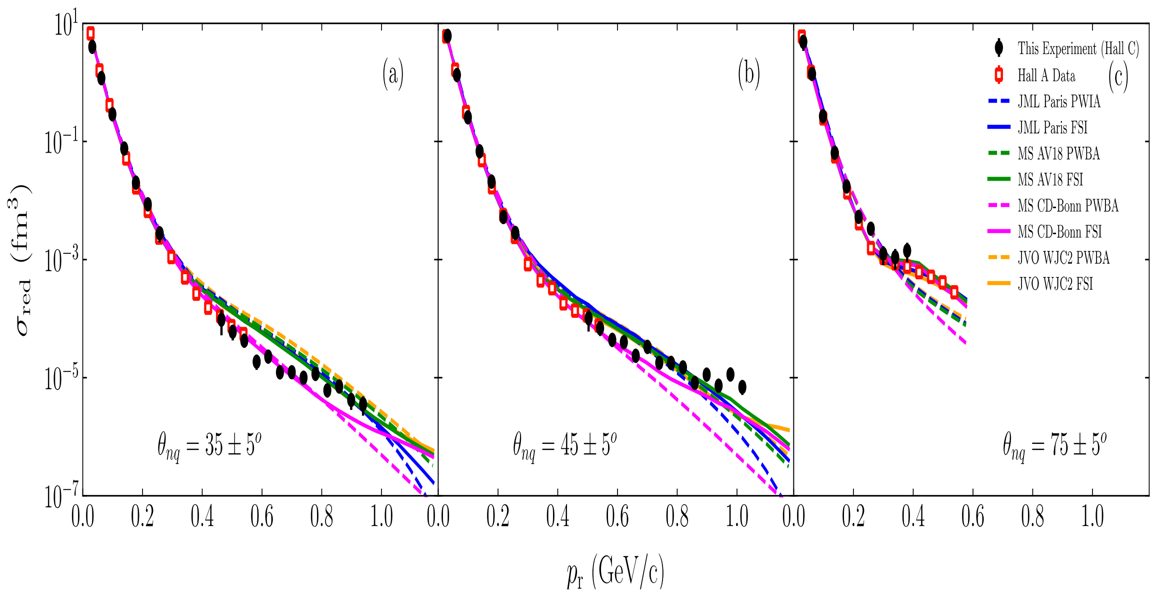
$$|p_i| = |p_f - q| \gtrsim 800\text{MeV/c}$$



3-days of commissioning measurement,



JLab experiment $Q^2 = 4\ GeV^2$

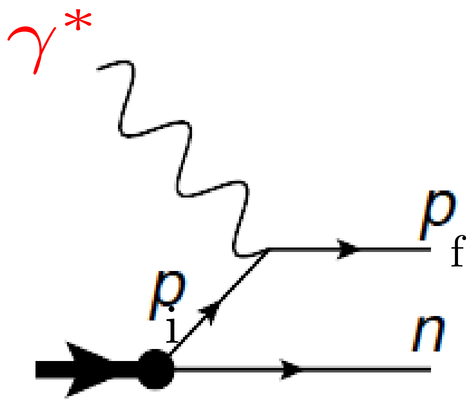


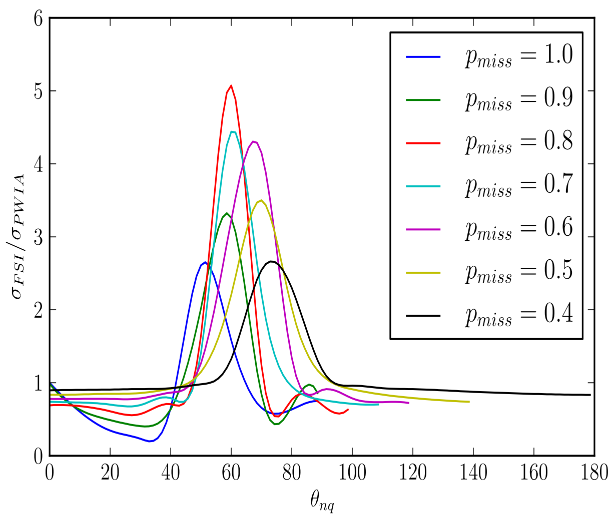
C. Yero, et al Phys.Rev.Lett. 125 (2020) 26, 262501

Possibility of Experimental Verification

Considering reaction: $e + d \to e' + p_f + n$

$$|p_i| = |p_f - q| \gtrsim 800\text{MeV}/c$$





PAC-49, 2021

PAC 49 SUMMARY OF JEOPARDY RECOMMENDATIONS							
Number	Contact Person	Title	Hall	Previously Approved Days	Days Already Rec'd	Days Awarded	PAC Decision
E12-09-011	Tanja Horn	Studies of the L-T Separated Kaon Electroproduction Cross Section from 5-11 GeV	C	40	32	8	Remain active
E12-10-003	W. Boeglin	Deuteron Electro-Disintegration at Very High Missing Momentum	C	21	3	18	Upgrade Rating to A-

1) Is there any new information that would affect the scientific importance or impact of the Experiment since it was originally proposed?

PAC 36 graded the proposal with B+ because, even though the physics motivation was viewed highly, the foreseen impact of the result was judged to be limited. The results of the three days commissioning in April 2018, published in Physical Review Letters 125, 262501 (2020), exhibit an unexpected behavior when compared with theoretical calculations. Therefore, the expected impact of future data has increased.

Outlook on Experimental Verification of the Effect

- analysis of the experiment will require careful account for competing nuclear effects most importantly final state interactions

- If angular dependence is found it will motivate new area of research
  a: modeling non-nucleonic components in the deuteron,
  b: understanding their origin and nature
  c: evaluating parameters that can be used for Equation of State of high density Nuclear Matter

- If no angular dependence is found,
  a: nucleonic degrees persist at very high density fluctuations
  b: non-nucleonic components conspire to preserve angular condition
  c: theory was wrong

Second Property: Inverse Fractional Dependence of High Momentum Component

$$a_{NN}(A,y) \approx a_{NN}(A,0) \cdot f(y)$$ with $f(0) = 1$ and $f(1) = 0$

$$f(|x_p - x_n|) = 1 - \sum_{j=1}^{n} b_i |x_p - x_x|^i \text{ with } \sum_{j=1}^{n} b_i = 0$$

In the limit $\sum_{j=1}^{n} b_i |x_p - x_x|^i \ll 1$ Momentum distributions of p & n are inverse proportional to their fractions

$$n^A_{p/n}(p) \approx \frac{1}{2x_{p/n}} a_2(A,y) \cdot n_d(p)$$